

Lecture 18: Continuum Thermo & Lagrangian Bal.

Logistics: - PS 7 is due

- probably no new PS this week

Last time: Eulerian balance laws

$$\frac{d}{dt} \int_{\Omega_t} \rho \phi \, dV_x = \int_{\Omega_t} \rho \dot{\phi} \, dV_x$$

$$\text{mass: } \dot{\rho} + \rho \nabla_x \cdot \underline{v} = 0$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \underline{v}) = 0$$

$$\text{lin. mom: } \rho \dot{\underline{v}} = \nabla \cdot \underline{\underline{\sigma}} + \rho \underline{b}$$

$$\frac{\partial}{\partial t} (\rho \underline{v}) + \nabla \cdot (\rho \underline{v} \otimes \underline{v} - \underline{\underline{\sigma}}) = \rho \underline{b}$$

$$\text{ang. mom: } \underline{\underline{\sigma}}^T = \underline{\underline{\sigma}}$$

rate of net working:

$$\frac{d}{dt} K[\Omega_t] = \int_{\Omega_t} \rho \underline{v} \cdot \dot{\underline{v}} \, dV$$

$$P[\Omega_t] = \int_{\Omega_t} \rho \underline{b} \cdot \underline{v} \, dV + \int_{\partial \Omega_t} \underline{t} \cdot \underline{v} \, dA_x$$

$$W[\Omega_t] = \int_{\Omega_t} \underline{\underline{\sigma}} : \underline{\underline{d}} \, dV_x$$

Today: Eulerian Energy & Entropy balance

Lagrangian balance laws

Local Eulerian form of First law of Thermo.

Integral form

$$\frac{d}{dt} U[\Omega_t] = Q[\Omega_t] + W[\Omega_t]$$

where $U[\Omega_t] = \int_{\Omega_t} \rho \phi dV_x$

$$Q[\Omega_t] = \int_{\Omega_t} \rho r dV_x - \int_{\partial\Omega_t} \mathbf{q} \cdot \underline{n} dA_x$$

$$W[\Omega_t] = \int_{\Omega_t} \underline{\underline{\sigma}} : \underline{\underline{d}} dV_x$$

Substituting

$$\frac{d}{dt} \int_{\Omega_t} \rho \phi dV_x = \int_{\Omega_t} \underline{\underline{\sigma}} : \underline{\underline{d}} dV_x - \int_{\partial\Omega_t} \mathbf{q} \cdot \underline{n} dA_x + \int_{\Omega_t} \rho r dV_x$$

using deriv. rel. to mass & div. theorem

$$\int_{\Omega_t} \rho \dot{\phi} - \underline{\underline{\sigma}} : \underline{\underline{d}} + \nabla_x \cdot \mathbf{q} - \rho r dV_x = 0$$

by arbitrary vol. of Ω_t

$$\rho \dot{\phi} = \underline{\underline{\sigma}} : \underline{\underline{d}} - \nabla_x \cdot \mathbf{q} + \rho r$$

local eulerian
form of energy bal.

for cons. form expand $\dot{\phi} = \frac{\partial \phi}{\partial t} + \underline{v} \cdot \nabla \phi$

$$\rho \dot{\phi} = \rho \frac{\partial \phi}{\partial t} + \rho \underline{v} \cdot \nabla_x \phi = \frac{\partial}{\partial t} (\rho \phi) - \phi \frac{\partial \rho}{\partial t} + \rho \underline{v} \cdot \nabla \phi$$

mass balance: $\frac{\partial \rho}{\partial t} = -\nabla_x \cdot (\rho \underline{v})$

$$\begin{aligned} \rho \dot{\phi} &= \frac{\partial}{\partial t} (\rho \phi) + \phi \nabla_x \cdot (\rho \underline{v}) + \rho \underline{v} \cdot \nabla \phi \\ &= \frac{\partial}{\partial t} (\rho \phi) + \nabla_x \cdot (\rho \phi \underline{v}) \end{aligned}$$

→ sub. conservative form

$$\boxed{\frac{\partial}{\partial t} (\rho \phi) + \nabla_x \cdot [\rho \phi \underline{v} + \underline{q}] = \underline{\sigma} : \underline{d} + \rho r}$$

Local Eulerian form of Second Law

Integral form of Clausius-Duhem form of 2nd law

$$\frac{d}{dt} \int_{\Omega_t} \rho s \, dV_x \geq \int_{\Omega_t} \frac{\rho r}{\theta} \, dV_x - \int_{\partial \Omega_t} \frac{\underline{q} \cdot \underline{n}}{\theta} \, dA_x$$

using deriv. with respect to mass + div thm

$$\int_{\Omega_t} \rho \dot{s} \, dV_x \geq \int_{\Omega} \frac{\rho r}{\theta} - \nabla_x \cdot \frac{\underline{q}}{\theta} \, dV_x$$

→ localizing

$$\boxed{\rho \dot{s} \geq \frac{\rho r}{\theta} - \nabla_x \cdot \left(\frac{\underline{q}}{\theta} \right)}$$

Clausius-Duhem
in local Eulerian form

$$\begin{aligned} \theta \rho \dot{s} &\geq \rho r - \theta \left(\frac{1}{\theta} \nabla_x \cdot \mathbf{q} + \mathbf{q} \cdot \nabla_x \left(\frac{1}{\theta} \right) \right) \\ &\geq \rho r - \nabla_x \cdot \mathbf{q} + \frac{1}{\theta} \nabla_x \theta \end{aligned}$$

introduce internal dissipation density

$$\delta = \theta \rho \dot{s} - (\rho r - \nabla_x \cdot \mathbf{q})$$

difference between entropy increase

and local rate of heating

$$dS \geq \frac{dQ}{\theta}$$

$$\delta - \frac{1}{\theta} \mathbf{q} \cdot \nabla_x \theta \geq 0$$

I, Any point where $\nabla_x \theta = 0 \Rightarrow \delta \geq 0$

\Rightarrow bodies with uniform θ have non. neg. dissip.

II If $\delta = 0$, i.e. reversible process

$$\Rightarrow \mathbf{q} \cdot \nabla_x \theta \leq 0$$

angle between \mathbf{q} and $\nabla_x \theta > 90$



q = energy density

⇒ heat flows down temp. gradient

$$\text{Fourier's law: } \underline{q} = -\underline{\kappa} \nabla \theta$$

To study the consequences of Clausius-Duhem for constitutive laws we introduce

Helmholtz free energy density:

$$\psi(\underline{x}, t) = \phi(\underline{x}, t) - \theta(\underline{x}, t) s(\underline{x}, t)$$

$$H = U - TS$$

Reformulate CD for in terms of ψ

Material derivative

$$\begin{aligned} \frac{d}{dt}(\theta s) &= \frac{\partial}{\partial t}(\theta s) + \nabla_{\underline{x}}(\theta s) \cdot \underline{v} \\ &= \theta \frac{\partial s}{\partial t} + s \frac{\partial \theta}{\partial t} + \theta \nabla_{\underline{x}} s \cdot \underline{v} + s \nabla_{\underline{x}} \theta \cdot \underline{v} \\ &= \theta \left(\frac{\partial s}{\partial t} + \nabla_{\underline{x}} s \cdot \underline{v} \right) + s \left(\frac{\partial \theta}{\partial t} + \nabla_{\underline{x}} \theta \cdot \underline{v} \right) \\ &= \theta \dot{s} + s \dot{\theta} \end{aligned}$$

from def. $\psi = \phi - \theta s$

$$\dot{\psi} = \dot{\phi} - \theta \dot{s} - s \dot{\theta} \Rightarrow \dot{\phi} = \dot{\psi} + \theta \dot{s} + s \dot{\theta}$$

subst into energy cons.

$$\rho \dot{\phi} = \underline{\underline{\sigma}} : \underline{\underline{d}} - \nabla_x \cdot \underline{q} + \rho r$$

$$\rho \dot{\psi} + \rho \theta \dot{s} + \rho s \dot{\theta} =$$

$$\rho \theta \dot{s} = \underline{\underline{\sigma}} : \underline{\underline{d}} - \nabla_x \cdot \underline{q} + \rho r - \rho \dot{\psi} - \rho s \dot{\theta}$$

substitute into 2nd law

$$\rho \theta \dot{s} \geq \rho r - \nabla_x \cdot \underline{q} + \frac{1}{\theta} \underline{q} \cdot \nabla_x \theta$$

$$\underline{\underline{\sigma}} : \underline{\underline{d}} - \cancel{\nabla_x \cdot \underline{q}} + \cancel{\rho r} - \rho \dot{\psi} - \rho s \dot{\theta} \geq \cancel{\rho r} - \cancel{\nabla_x \cdot \underline{q}} + \frac{1}{\theta} \underline{q} \cdot \nabla_x \theta$$

solve $\rho \dot{\psi}$

$$\rho \dot{\psi} \leq \underline{\underline{\sigma}} : \underline{\underline{d}} - \rho s \dot{\theta} - \frac{1}{\theta} \underline{q} \cdot \nabla_x \theta$$

reduced Clausius-Duhem inequality

- independent of \underline{q} and r

In a body with const. θ $\dot{\theta} = 0$ $\nabla_x \theta = 0$

$$\rho \dot{\psi} \leq \underline{\underline{\sigma}} : \underline{\underline{d}}$$

for a reversible process $\Rightarrow \rho \dot{\psi} = \underline{\underline{\sigma}} : \underline{\underline{d}}$

\Rightarrow rate of change of Helmholtz free energy

is equal to stress power in rev. process
 \Rightarrow Helmholtz free energy is the part
of the internal energy available
for performing work at const. θ .

Summary Eulerian Balance Laws

Governing equations:

mass: $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \underline{v}) = 0$ 1

lin. mom: $\frac{\partial (\rho \underline{v})}{\partial t} + \nabla \cdot (\rho \underline{v} \otimes \underline{v} - \underline{\underline{\sigma}}) = \rho \underline{b}$ 3

ang. mom: $\underline{\underline{\sigma}} = \underline{\underline{\sigma}}^T$ 3

energy: $\frac{\partial (\rho \phi)}{\partial t} + \nabla \cdot (\rho \phi \underline{v} + \underline{q}) = \underline{\underline{\sigma}} : \underline{d} + \rho r$ 1

kinematic: $\underline{v}_m = \frac{\partial \varphi}{\partial t}$ + 3
11

Unknown fields

φ	\underline{v}	ρ	$\underline{\underline{\sigma}}$	θ	\underline{q}	ϕ	
3	3	1	9	1	3	1	= 21

We have 21 unknown fields but only 11 equations.

⇒ under constrained

⇒ Need additional constitutive eqns

Remarks:

1) Eulerian formulation of balance laws is independent of φ !

The motion is only needed to determine shape of domain. If \mathcal{B}_E is known we have 8 eqns and 18 unknowns

Need constitutive eqns that relate

$\underline{\epsilon}$, η , ϕ (secondary) to ρ , \underline{v} and θ (primary unknowns)

2) If thermal effects are neglected:

η , ϕ & θ disappears reducing the unknowns

to 12. The number of eqns reduces

to 7. Need constitutive eqns
that relate $\underline{\sigma}$ to ρ and \underline{v} .