

Lecture 21: Constitutive Theory

Logistics: - PS 8 is due Thursday

Last time: - Frame indifference

- two observers φ, φ^*

- Frame indifference

$$\phi^*(\underline{x}^*, t) = \phi(\underline{x}, t)$$

$$\underline{w}^*(\underline{x}^*, t) = \underline{Q} \underline{w}(\underline{x}, t)$$

$$\underline{s}^*(\underline{x}^*, t) = \underline{Q} \underline{s}(\underline{x}, t) \underline{Q}^T$$

- Velocity & Acceleration are not frame indifferent

\Rightarrow Centrifugal & Coriolis acc.

- Functions of fields

$$\hat{\underline{s}}(\underline{Q} \underline{s} \underline{Q}^T) = \underline{Q} \hat{\underline{s}}(\underline{s}) \underline{Q}^T$$

Today: - Isotropic tensor functions

- Fourth order tensors

- Material constraints

Isotropic Functions

Functions that are frame indifferent are also called isotropic. For two ref. frames related by rotation \underline{Q} we have

$$\begin{aligned}
 \phi(\theta(\underline{x}^*, t)) &= \phi(\theta(\underline{x}, t)) & \phi(\underline{Q}\underline{v}) &= \phi(\underline{v}) & \phi(\underline{Q}\underline{S}\underline{Q}^T) &= \phi(\underline{S}) \\
 \underline{u}(\theta(\underline{x}^*)) &= \underline{Q}\underline{u}(\theta(\underline{x})) & \underline{u}(\underline{Q}\underline{v}(\underline{x}^*)) &= \underline{Q}\underline{u}(\underline{v}(\underline{x})) & \underline{u}(\underline{Q}\underline{S}\underline{Q}^T) &= \underline{Q}\underline{u}(\underline{S}) \\
 \underline{\sigma}(\theta(\underline{x}^*)) &= \underline{Q}\underline{\sigma}(\theta(\underline{x}))\underline{Q}^T & \underline{\sigma}(\underline{Q}\underline{v}(\underline{x}^*)) &= \underline{Q}\underline{\sigma}(\underline{v}(\underline{x}))\underline{Q}^T & \underline{\sigma}(\underline{Q}\underline{S}\underline{Q}^T) &= \underline{Q}\underline{\sigma}\underline{Q}^T
 \end{aligned}$$

ϕ = scalar val. fun \underline{u} = vector val. fun $\underline{\sigma}$ = tensor val. fun

θ = scalar field \underline{v} = vector field \underline{S} = tensor field

Example: 1) $\phi(\underline{S}) = \det(\underline{S})$

$$\begin{aligned}
 \phi(\underline{Q}\underline{S}\underline{Q}^T) &= \det(\underline{Q}) \det(\underline{S}) \det(\underline{Q}^T) \\
 &= \phi(\underline{S})
 \end{aligned}$$

2) $\underline{u}(\underline{v}, \underline{A}) = \underline{A}\underline{v}$

$$\begin{aligned}
 \underline{u}(\underline{Q}\underline{v}, \underline{Q}\underline{A}\underline{Q}^T) &= \underline{Q}\underline{A}\underline{Q}^T \underline{Q}\underline{v} \\
 &= \underline{Q}\underline{u}(\underline{v}, \underline{A})
 \end{aligned}$$

Representation of isotropic tensor functions

The most general form of an isotropic tensor function that maps one symmetric tensor to another sym. tensor is given by

$$\underline{\underline{G}}(\underline{\underline{A}}) = \alpha_0(\mathcal{I}_A) \underline{\underline{I}} + \alpha_1(\mathcal{I}_A) \underline{\underline{A}} + \alpha_2(\mathcal{I}_A) \underline{\underline{A}}^2$$

Rivlin-Ericksen
Representation
Theorem

whr $\alpha_0, \alpha_1, \alpha_2$ are functions of

The set of principal invariants of $\underline{\underline{A}}$

- $\underline{\underline{G}}$ is clearly sym. if $\underline{\underline{A}}$ is sym.
- To see that $\underline{\underline{G}}$ is isotropic

$$\underline{\underline{G}}(\underline{\underline{Q}} \underline{\underline{A}} \underline{\underline{Q}}^T) = \alpha_0 \underline{\underline{I}} + \alpha_1 \underline{\underline{Q}} \underline{\underline{A}} \underline{\underline{Q}}^T + \alpha_2 \underbrace{\underline{\underline{Q}} \underline{\underline{A}} \underline{\underline{Q}}^T \underline{\underline{Q}} \underline{\underline{A}} \underline{\underline{Q}}^T}_{\underline{\underline{Q}} \underline{\underline{A}}^2 \underline{\underline{Q}}^T}$$

$$\underline{\underline{Q}} \underline{\underline{G}}(\underline{\underline{A}}) \underline{\underline{Q}}^T = \alpha_0 \underline{\underline{Q}} \underline{\underline{I}} \underline{\underline{Q}}^T + \alpha_1 \underline{\underline{Q}} \underline{\underline{A}} \underline{\underline{Q}}^T + \alpha_2 \underline{\underline{Q}} \underline{\underline{A}}^2 \underline{\underline{Q}}^T \quad \checkmark$$

because coeff. only depend on invariants of $\underline{\underline{A}}$

$\Rightarrow \underline{\underline{G}}(\underline{\underline{A}})$ isotropic!

A second representation can be obtained by eliminating $\underline{\underline{A}}^2$ term using Cayley-Hamilton

$$\underline{\underline{A}}^3 - I_1(\underline{\underline{A}}) \underline{\underline{A}}^2 + I_2(\underline{\underline{A}}) \underline{\underline{A}} - I_3(\underline{\underline{A}}) \underline{\underline{I}} = 0$$

multiply $\underline{\underline{A}}^{-1}$

$$\underline{\underline{A}}^2 - I_1 \underline{\underline{A}} + I_2 \underline{\underline{I}} - I_3 \underline{\underline{A}}^{-1} = 0$$

$$\underline{\underline{A}}^2 = I_1 \underline{\underline{A}} - I_2 \underline{\underline{I}} + I_3 \underline{\underline{A}}^{-1}$$

substitute into $G(\underline{\underline{A}})$

$$\boxed{G(\underline{\underline{A}}) = \beta_0(\underline{\underline{I}}) \underline{\underline{I}} + \beta_1(\underline{\underline{I}}) \underline{\underline{A}} + \beta_2(\underline{\underline{I}}) \underline{\underline{A}}^{-1}}$$

$$\beta_0 = \alpha_0 - I_2 \alpha_2 \quad \beta_1 = \alpha_1 - I_1 \alpha_2 \quad \beta_2 = I_3 \alpha_2$$

⇒ used in hyperelastic materials.

Isotropic 4-th order tensors

If $\underline{\underline{G}}(\underline{\underline{A}})$ is linear then it can be written

$$\underline{\underline{G}}(\underline{\underline{A}}) = \mathbb{C} \underline{\underline{A}}$$

where \mathbb{C} is a 4-th order tensor.

If we also ,