

## Second-order Tensors

Here we are interested in second-order tensors

Linear operators :  $\underline{v} = \underline{A} \underline{u}$

maps vector  $\underline{u} \in \mathcal{V}$  into vector  $\underline{v} \in \mathcal{V}$

Linearity requires that

$$1) \quad \underline{A}(\underline{u} + \underline{v}) = \underline{A}\underline{u} + \underline{A}\underline{v} \quad \text{for all } \underline{u}, \underline{v} \in \mathcal{V}$$

$$2) \quad \underline{A}(\alpha \underline{v}) = \alpha \underline{A}\underline{v} \quad \text{for all } \alpha \in \mathbb{R} \text{ and } \underline{v} \in \mathcal{V}$$

Example:  $\underline{A}$  maps every  $\underline{v} \in \mathcal{V}$  into  $\underline{n} \neq \underline{0} \in \mathcal{V}$ .

Is  $\underline{A}$  a tensor?

Consider  $\underline{u}, \underline{v}, \underline{w} \in \mathcal{V}$

$$\underline{w} = \underline{u} + \underline{v}$$

$$\underline{A}(\underline{u} + \underline{v}) \stackrel{?}{=} \underline{A}\underline{u} + \underline{A}\underline{v}$$

$$\underline{A}\underline{w} \stackrel{?}{=} \underline{A}\underline{u} + \underline{A}\underline{v}$$

$$\underline{n} \neq \underline{n} + \underline{n}$$

$\Rightarrow \underline{A}$  is not a tensor, because it is not linear

## Tensor algebra

For all  $\underline{v} \in \mathcal{V}$  we define

$$1) \quad (\underline{\alpha A}) \underline{v} = \underline{A} (\underline{\alpha v}) \quad \text{scalar multiplication}$$

$$2) \quad (\underline{A} + \underline{B}) \underline{v} = \underline{A} \underline{v} + \underline{B} \underline{v} \quad \text{tensor sum}$$

$$3) \quad (\underline{A} \underline{B}) \underline{v} = \underline{A} (\underline{B} \underline{v}) \quad \text{tensor product}$$

Note there is also a scalar product introduced later.

The set of all  $2^{\text{nd}}$ -order tensors  $\mathcal{V}^2$  is  
a vector space

$$1) \quad \alpha \underline{A} \in \mathcal{V}^2 \quad \text{for all } \underline{A} \in \mathcal{V} \text{ and } \alpha \in \mathbb{R}$$

$$2) \quad \underline{A} + \underline{B} \in \mathcal{V}^2 \quad \text{for all } \underline{A}, \underline{B} \in \mathcal{V}^2$$

$$3) \quad \underline{A} \underline{B} \in \mathcal{V}^2 \quad \text{for all } \underline{A}, \underline{B} \in \mathcal{V}^2$$

Any of these operations will produce  
another second-order tensor.

Q: What is a basis for  $\mathcal{V}^2$ ?

Two tensors  $\underline{\underline{A}}$  and  $\underline{\underline{B}}$  are equal if

$$\underline{\underline{A}} \underline{v} = \underline{\underline{B}} \underline{v} \quad \text{for all } \underline{v} \in \mathcal{V}$$

Zero tensor:  $\underline{\underline{0}} \underline{v} = \underline{0}$  for all  $\underline{v} \in \mathcal{V}$

Identity tensor:  $\underline{\underline{I}} \underline{v} = \underline{v}$  for all  $\underline{v} \in \mathcal{V}$

## Representation of a tensor

In a frame  $\{\underline{e}_i\}$  a second order tensor

$\underline{\underline{S}}$  is represented by nine numbers

$$S_{ij} = \underline{e}_i \cdot \underline{\underline{S}} \underline{e}_j$$

Matrix representation of tensor in  $\{\underline{e}_i\}$

$$[\underline{\underline{S}}] = \begin{bmatrix} S_{11} & S_{12} & S_{13} \\ S_{21} & S_{22} & S_{23} \\ S_{31} & S_{32} & S_{33} \end{bmatrix} \in \mathbb{R}^3 \times \mathbb{R}^3$$

Note that  $[\underline{\underline{S}}]_{ij} = S_{ij}$

Consider  $\underline{v} = \underline{S} \underline{u}$  where  $\underline{v} = v_k \underline{e}_k$ ,  $\underline{u} = u_j \underline{e}_j$

$$v_k \underline{e}_k = \underline{S} (u_j \underline{e}_j) = \underline{S} \underline{e}_j u_j$$

multiply by  $\underline{e}_i$  from left

$$v_k \underline{e}_i \cdot \underline{e}_k = \underline{e}_i \cdot \underline{S} \underline{e}_j u_j$$

$$v_k \delta_{ik} = \underline{e}_i \cdot \underline{S} \underline{e}_j u_j$$

$$v_i = (\underline{e}_i \cdot \underline{S} \underline{e}_j) u_j$$

$$v_i = S_{ij} u_j$$

## Dyadic Product

The dyadic product of two vectors  $\underline{a}$  and  $\underline{b}$  is the 2<sup>nd</sup>-order tensor  $\underline{a} \otimes \underline{b}$  defined by

$$(\underline{a} \otimes \underline{b}) \underline{v} = (\underline{b} \cdot \underline{v}) \underline{a} \quad \text{for all } \underline{v} \in \mathcal{V}$$

This has the form:  $\underline{A} \underline{v} = \alpha \underline{a}$

in components:  $A_{ij} v_j = \alpha a_i$

$$\alpha = \underline{b} \cdot \underline{v} = b_j v_j$$

$$A_{ij} = [\underline{a} \otimes \underline{b}]_{ij}$$

$$\Rightarrow [\underline{a} \otimes \underline{b}]_{ij} v_j = b_j v_j a_i$$

$$[\underline{a} \otimes \underline{b}]_{ij} v_j = (a_i b_j) v_j$$

$$\Rightarrow [\underline{a} \otimes \underline{b}]_{ij} = a_i b_j$$

So that

$$[\underline{a} \otimes \underline{b}] = \begin{bmatrix} a_1 b_1 & a_1 b_2 & a_1 b_3 \\ a_2 b_1 & a_2 b_2 & a_2 b_3 \\ a_3 b_1 & a_3 b_2 & a_3 b_3 \end{bmatrix} = \underline{a} \underline{b}^T$$

Linearity of dyadic product:

for scalars  $\alpha, \beta \in \mathbb{R}$  and vectors  $\underline{a}, \underline{b}, \underline{v}, \underline{w} \in V$

$$(\underline{a} \otimes \underline{b})(\alpha \underline{v} + \beta \underline{w}) = \alpha (\underline{a} \otimes \underline{b}) \underline{v} + \beta (\underline{a} \otimes \underline{b}) \underline{w}$$

The product of two dyadic products

$$(\underline{a} \otimes \underline{b})(\underline{c} \otimes \underline{d}) = (\underline{b} \cdot \underline{c}) \underline{a} \otimes \underline{d} \Rightarrow \text{HW2}$$

needed for tensor product.

## Basis for $V^2$

Given any frame  $\{\underline{e}_i\}$  the nine dyadic products  $\{\underline{e}_i \otimes \underline{e}_j\}$  form a basis for  $V^2$ .

Any second-order tensor  $\underline{\underline{S}}$  can be written as linear combination

$$\underline{\underline{S}} = S_{ij} \underline{e}_i \otimes \underline{e}_j$$

where  $S_{ij} = \underline{e}_i \cdot \underline{\underline{S}} \underline{e}_j$

Consider  $\underline{v} = \underline{\underline{S}} \underline{u}$  with  $\underline{v} = v_i \underline{e}_i$ ,  $\underline{u} = u_k \underline{e}_k$

$$v_i \underline{e}_i = S_{ij} (\underline{e}_i \otimes \underline{e}_j) (u_k \underline{e}_k)$$

$$= S_{ij} u_k (\underline{e}_i \otimes \underline{e}_j) \cdot \underline{e}_k \quad \text{apply def. of dyadic}$$

$$= S_{ij} u_k (\underline{e}_j \cdot \underline{e}_k) \underline{e}_i = S_{ij} u_k \delta_{kj} \underline{e}_i$$

$$v_i \underline{e}_i = S_{ij} u_j \underline{e}_i$$

$$v_i = S_{ij} u_j$$

used often

## Tensor algebra in components

Addition:  $\underline{\underline{H}} = \underline{\underline{S}} + \underline{\underline{T}}$

$$H_{ij} \underline{e}_i \otimes \underline{e}_j = S_{ij} \underline{e}_i \otimes \underline{e}_j + T_{ij} \underline{e}_i \otimes \underline{e}_j \\ = (S_{ij} + T_{ij}) \underline{e}_i \otimes \underline{e}_j$$

$$\boxed{H_{ij} = S_{ij} + T_{ij}}$$

Scalar multiplication:  $\underline{\underline{H}} = \alpha \underline{\underline{S}} \Rightarrow \boxed{H_{ij} = \alpha S_{ij}}$

Product:  $\underline{\underline{H}} = \underline{\underline{S}} \underline{\underline{T}}$

$$\underline{\underline{H}} = S_{ij} (\underline{e}_i \otimes \underline{e}_j) T_{kl} (\underline{e}_k \otimes \underline{e}_l) \\ = S_{ij} T_{kl} \underbrace{(\underline{e}_i \otimes \underline{e}_j)(\underline{e}_k \otimes \underline{e}_l)}_{\text{product of two dyads}}$$

$$= S_{ij} T_{kl} (\underline{e}_j \otimes \underline{e}_k) \underline{e}_i \otimes \underline{e}_l \\ \delta_{jk}$$

$$= S_{ij} T_{jl} \underline{e}_i \otimes \underline{e}_l$$

$$H_{il} \underline{e}_i \otimes \underline{e}_l = S_{ij} T_{jl} \underline{e}_i \otimes \underline{e}_l$$

$$\Rightarrow \boxed{H_{il} = S_{ij} T_{jl}} \quad \text{note the dummy } j!$$

## Transpose of a tensor

To any  $\underline{\underline{S}} \in \mathcal{V}^2$  we associate a transpose  $\underline{\underline{S}}^T \in \mathcal{V}^2$  the unique tensor such that

$$\underline{\underline{S}} \underline{\underline{u}} \cdot \underline{\underline{v}} = \underline{\underline{u}} \cdot \underline{\underline{S}}^T \underline{\underline{v}} \quad \text{for all } \underline{\underline{u}}, \underline{\underline{v}} \in \mathcal{V}$$

This implies that  $S_{ij}^T = S_{ji}$  as follows

$$(S_{ij} u_j e_i) \cdot (v_l e_l) = (u_k e_k) \cdot (S_{ij}^T v_j e_i)$$

$$S_{ij} u_j v_l (e_i \cdot e_l) = S_{ij}^T v_j u_k (e_k \cdot e_i)$$

$$S_{ij} u_j v_l \delta_{il} = S_{ij}^T v_j u_k \delta_{ki}$$

$$S_{ij} u_j v_i = S_{ij}^T v_j u_i$$

rename indices  
 $i \leftrightarrow j$  on rhs

$$S_{ij} u_j v_i = S_{ji}^T u_j v_i$$

$$\Rightarrow S_{ij} = S_{ji}^T \quad \checkmark$$

Properties of transpose:

$$(\underline{\underline{A}}^T)^T = \underline{\underline{A}}$$

$$(\underline{\underline{A}} \underline{\underline{B}})^T = \underline{\underline{B}}^T \underline{\underline{A}}^T$$

$$(\underline{\underline{u}} \otimes \underline{\underline{v}})^T = \underline{\underline{v}} \otimes \underline{\underline{u}}$$



$\underline{\underline{S}}$  is symmetric if  $\underline{\underline{S}} = \underline{\underline{S}}^T$   $S_{ij} = S_{ji}$   
 $\underline{\underline{S}}$  is skew-symmetric if  $\underline{\underline{S}} = -\underline{\underline{S}}^T$   $S_{ij} = -S_{ji}$

Symmetric - Skew decomposition:  
 Any tensor  $\underline{\underline{S}} \in \mathcal{V}^2$  can be written as

$$\begin{aligned}
 \underline{\underline{S}} &= \underline{\underline{E}} + \underline{\underline{W}} \\
 \underline{\underline{E}} &= \frac{1}{2} (\underline{\underline{S}} + \underline{\underline{S}}^T) & \underline{\underline{E}} &= \underline{\underline{E}}^T \\
 \underline{\underline{W}} &= \frac{1}{2} (\underline{\underline{S}} - \underline{\underline{S}}^T) & \underline{\underline{W}} &= -\underline{\underline{W}}^T
 \end{aligned}$$

Note:  $\underline{\underline{W}} = \begin{bmatrix} 0 & W_{12} & W_{13} \\ -W_{12} & 0 & W_{23} \\ -W_{13} & -W_{23} & 0 \end{bmatrix}$  only 3 indep. comp.

$\Rightarrow$  can be related to an axial vector  $\underline{\underline{w}}$

$$\underline{\underline{W}} \underline{\underline{v}} = \underline{\underline{w}} \times \underline{\underline{v}} \quad \text{for all } \underline{\underline{v}} \in \mathcal{V}$$

Relation:

$$\begin{aligned}
 W_{ij} &= -\epsilon_{ijk} w_k \\
 w_k &= -\frac{1}{2} \epsilon_{ijk} W_{ij}
 \end{aligned}
 \quad
 \underline{\underline{W}} = \begin{bmatrix} 0 & -w_3 & w_2 \\ w_3 & 0 & -w_1 \\ -w_2 & w_1 & 0 \end{bmatrix}$$

## Trace of a tensor

We define the trace of a dyad as

$$\text{tr}(\underline{a} \otimes \underline{b}) = \underline{a} \cdot \underline{b} = a_i b_i$$

this implies that

$$\text{tr}(\underline{A}) = A_{ii} = A_{11} + A_{22} + A_{33}$$

as follows  $\text{tr}(A_{ij} \underline{e}_i \otimes \underline{e}_j) = A_{ij} \text{tr}(\underline{e}_i \otimes \underline{e}_j)$   
 $= A_{ij} \delta_{ij} = A_{ii}$

Properties:  $\text{tr}(\underline{A}^T) = \text{tr}(\underline{A})$

$$\text{tr}(\underline{A}\underline{B}) = \text{tr}(\underline{B}\underline{A})$$

$$\text{tr}(\underline{A} + \underline{B}) = \text{tr}(\underline{A}) + \text{tr}(\underline{B})$$

$$\text{tr}(\alpha \underline{A}) = \alpha \text{tr}(\underline{A})$$

Decomposition:  $\underline{A} = \alpha \underline{I} + \text{dev } \underline{A}$

Spherical tensor:  $\alpha \underline{I}$  where  $\alpha = \frac{1}{3} \text{tr}(\underline{A})$

Deviatoric tensor:  $\text{dev } \underline{A} = \underline{A} - \alpha \underline{I}$

$$\text{tr}(\text{dev } \underline{A}) = 0$$

## Tensor scalar product (Contraction)

analogous to scalar product of vectors

$$\underline{\underline{A}} : \underline{\underline{B}} = \text{tr}(\underline{\underline{A}}^T \underline{\underline{B}}) = A_{ij} B_{ij} \quad \text{scalar } \nabla$$

explicitly:

$$\begin{aligned} \underline{\underline{A}} : \underline{\underline{B}} &= \sum_{i=1}^3 \sum_{j=1}^3 A_{ij} B_{ij} = A_{11} B_{11} + A_{12} B_{12} + A_{13} B_{13} + \dots \\ &\quad A_{21} B_{21} + A_{22} B_{22} + A_{23} B_{23} + \dots \\ &\quad A_{31} B_{31} + A_{32} B_{32} + A_{33} B_{33} \end{aligned}$$

The index expression is derived as follows

$$\underline{\underline{A}} : \underline{\underline{B}} = \text{tr}(\underline{\underline{A}}^T \underline{\underline{B}}) =$$

$$\underline{\underline{A}}^T \underline{\underline{B}} = (A_{ji} \underline{e}_i \otimes \underline{e}_j) (B_{kl} \underline{e}_k \otimes \underline{e}_l)$$

$$= A_{ji} B_{kl} (\underline{e}_i \otimes \underline{e}_j) (\underline{e}_k \otimes \underline{e}_l)$$

$$= A_{ji} B_{kl} (\underline{e}_j \cdot \underline{e}_k) (\underline{e}_i \otimes \underline{e}_l) = A_{ji} B_{kl} \delta_{jk} \underline{e}_i \otimes \underline{e}_l$$

$$= A_{ji} B_{jl} \underline{e}_i \otimes \underline{e}_l$$

$$\text{tr}(\underline{\underline{A}}^T \underline{\underline{B}}) = A_{ji} B_{jl} \delta_{il} = A_{ji} B_{ji} = A_{ij} B_{ij} \quad \checkmark$$

Properties: 1)  $\underline{\underline{A}} : \underline{\underline{B}} = \underline{\underline{B}} : \underline{\underline{A}}$

$$2) (\underline{\underline{a}} \otimes \underline{\underline{b}}) : (\underline{\underline{c}} \otimes \underline{\underline{d}}) = (\underline{\underline{a}} \cdot \underline{\underline{c}})(\underline{\underline{b}} \cdot \underline{\underline{d}})$$

First follows from prop. of trace

$$\underline{\underline{A}} : \underline{\underline{B}} = \text{tr}(\underline{\underline{A}}^T \underline{\underline{B}}) = \text{tr}((\underline{\underline{A}}^T \underline{\underline{B}})^T) = \text{tr}(\underline{\underline{B}}^T \underline{\underline{A}}) = \underline{\underline{B}} : \underline{\underline{A}}$$

Second property  $[\underline{\underline{a}} \otimes \underline{\underline{b}}]_{ij} = a_i b_j$

$$\begin{aligned} (\underline{\underline{a}} \otimes \underline{\underline{b}}) : (\underline{\underline{c}} \otimes \underline{\underline{d}}) &= [\underline{\underline{a}} \otimes \underline{\underline{b}}]_{ij} [\underline{\underline{c}} \otimes \underline{\underline{d}}]_{ij} = a_i b_j c_i d_j \\ &= a_i c_i b_j d_j \\ &= (\underline{\underline{a}} \cdot \underline{\underline{c}})(\underline{\underline{b}} \cdot \underline{\underline{d}}) \end{aligned}$$

A common norm for tensors is

$$|\underline{\underline{A}}| = \sqrt{\underline{\underline{A}} : \underline{\underline{A}}^T} = \sqrt{A_{ij} A_{ij}} \geq 0$$

Note: Tensor scalar product will be important to express the work done during deformation.

For example the shear heating in glaciology.

## Determinant and Inverse

The determinant of  $\underline{\underline{A}} \in \mathcal{V}^2$  is the scalar

$$\det(\underline{\underline{A}}) = \det \begin{vmatrix} A_{11} & A_{12} & A_{13} \\ A_{21} & A_{22} & A_{23} \\ A_{31} & A_{32} & A_{33} \end{vmatrix} = \epsilon_{ijk} [\underline{\underline{A}}]_{i1} [\underline{\underline{A}}]_{j2} [\underline{\underline{A}}]_{k3}$$

where  $[\underline{\underline{A}}]_{i1}$ ,  $[\underline{\underline{A}}]_{j2}$ ,  $[\underline{\underline{A}}]_{k3}$  are the columns of  $[\underline{\underline{A}}]$

Properties:  $\det(\underline{\underline{A}}\underline{\underline{B}}) = \det(\underline{\underline{A}}) \det(\underline{\underline{B}})$

$$\det(\underline{\underline{A}}^T) = \det(\underline{\underline{A}})$$

$$\det(\alpha \underline{\underline{A}}) = \alpha^n \det(\underline{\underline{A}}) \quad (\underline{\underline{A}} \text{ is } n \times n)$$

$\underline{\underline{A}}$  is singular if  $\det \underline{\underline{A}} = 0$ .

If  $\det \underline{\underline{A}} \neq 0$  then the inverse  $\underline{\underline{A}}^{-1}$  exists

$$\underline{\underline{A}}^{-1} \underline{\underline{A}} = \underline{\underline{A}} \underline{\underline{A}}^{-1} = \underline{\underline{I}}$$

Properties:  $(\underline{\underline{A}} \underline{\underline{B}})^{-1} = \underline{\underline{B}}^{-1} \underline{\underline{A}}^{-1}$

$$(\underline{\underline{A}}^{-1})^{-1} = \underline{\underline{A}}$$

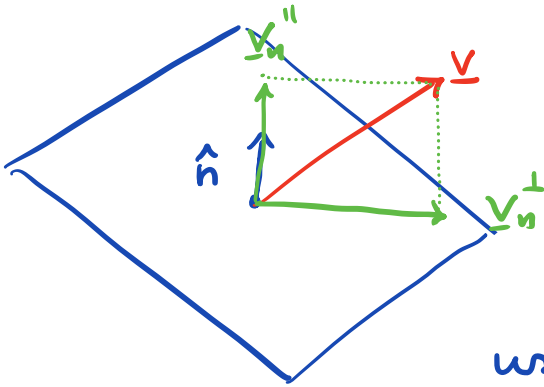
$$(\underline{\underline{A}}^{-1})^T = (\underline{\underline{A}}^T)^{-1}$$

$$(\alpha \underline{\underline{A}})^{-1} = \frac{1}{\alpha} \underline{\underline{A}}^{-1}$$

$$\det(\underline{\underline{A}}^{-1}) = \det(\underline{\underline{A}})^{-1} = \frac{1}{\det(\underline{\underline{A}})}$$

## Projection & Reflection tensors

commonly used to partition forces on a surface.



$$\underline{v} = \underline{v}_n^{\parallel} + \underline{v}_n^{\perp}$$

$$\underline{v}_n^{\parallel} = (\underline{v} \cdot \underline{\hat{n}}) \underline{\hat{n}}$$

$$\underline{v}_n^{\perp} = \underline{v} - \underline{v}_n^{\parallel}$$

use dyadic property

$$\underline{v}_n^{\parallel} = (\underline{v} \cdot \underline{\hat{n}}) \underline{\hat{n}} = (\underline{\hat{n}} \otimes \underline{\hat{n}}) \underline{v} = \underline{\underline{P}}_n^{\parallel} \underline{v}$$

$$\underline{v}_n^{\perp} = \underline{v} - (\underline{\hat{n}} \otimes \underline{\hat{n}}) \underline{v} = (\underline{\underline{I}} - \underline{\hat{n}} \otimes \underline{\hat{n}}) \underline{v} = \underline{\underline{P}}_n^{\perp} \underline{v}$$

$$\underline{\underline{P}}_n^{\parallel} = \underline{\hat{n}} \otimes \underline{\hat{n}}$$

$$\underline{\underline{P}}_n^{\perp} = \underline{\underline{I}} - \underline{\hat{n}} \otimes \underline{\hat{n}}$$

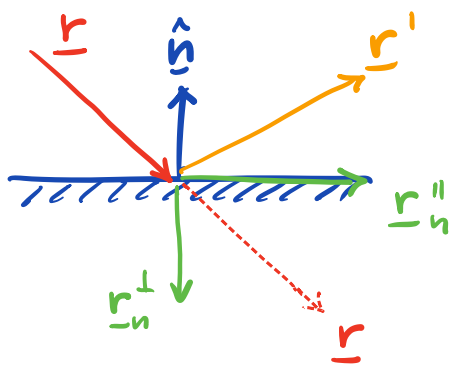
Properties:  $\underline{\underline{P}} = \underline{\underline{P}}^T$  symmetric

$$\underline{\underline{P}}^2 = \underline{\underline{P}}$$

$$\underline{\underline{P}} + \underline{\underline{P}}^t = \underline{\underline{I}}$$

$$\underline{\underline{P}} \underline{\underline{P}}^t = \underline{\underline{0}}$$

## Reflections



incoming:  $\underline{r} = \underline{r}_n^{\parallel} + \underline{r}_n^{\perp}$

reflected:  $\underline{r}' = \underline{r}_n^{\parallel} - \underline{r}_n^{\perp}$

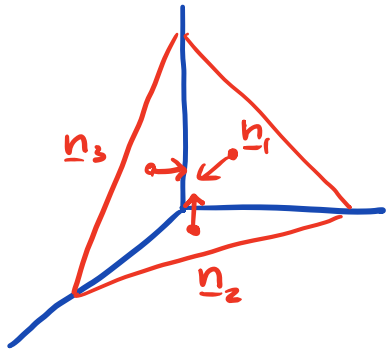
$$\underline{r}' = (\underline{\underline{P}}_n^{\parallel} - \underline{\underline{P}}_n^{\perp}) \underline{r}$$

$$\underline{r}' = (\underline{\underline{I}} - 2 \underline{\hat{n}} \otimes \underline{\hat{n}}) \underline{r}$$

$$\underline{r}' = \underline{\underline{R}}_n \underline{r}$$

Reflection tensor:  $\underline{\underline{R}}_n = \underline{\underline{I}} - 2 \underline{\hat{n}} \otimes \underline{\hat{n}}$

Example: Corner reflector



Inverts the direction of any ray that reflects off all three surfaces.

Direction of triply reflected ray:

$$\underline{r}''' = \underline{R}_{\underline{n}_1} \underline{R}_{\underline{n}_2} \underline{R}_{\underline{n}_3} \underline{r}$$

Show that  $\underline{r}''' = -\underline{r}$  !