

## Cauchy - Green Strain Tensor

Consider a deformation  $\varphi: \mathcal{B} \rightarrow \mathcal{B}'$  with  $\underline{\underline{F}} = \nabla \varphi$ , then the (right) Cauchy - Green strain tensor is

$$\underline{\underline{C}} = \underline{\underline{F}}^T \underline{\underline{F}}.$$

Note that  $\underline{\underline{C}}$  is always symmetric pos. definite.

The deformation gradient  $\underline{\underline{F}}$  contains information about both rotations and stretches. Using the right polar decomposition we have

$$\underline{\underline{F}} = \underline{\underline{R}} \underline{\underline{U}} \quad \begin{array}{l} \underline{\underline{R}} \text{ is rotation matrix} \\ \underline{\underline{U}} = \sqrt{\underline{\underline{F}}^T \underline{\underline{F}}} \text{ is right stretch tensor} \end{array}$$

Clearly  $\underline{\underline{C}} = \underline{\underline{U}}^2$  and the rotation  $\underline{\underline{R}}$  implicit in  $\underline{\underline{F}}$  is not present in  $\underline{\underline{C}}$ .

$\Rightarrow$  The right Cauchy Green strain tensor only contains information about stretches.

Hence we cannot obtain  $\underline{\underline{F}}$  from  $\underline{\underline{C}}$  !

## Remarks:

1) Strictly the right-stretch tensor  $\underline{\underline{U}}$  is sufficient.

We introduce  $\underline{\underline{C}} = \underline{\underline{U}}^2$  to avoid the tensor square root.

Simple example:

$$[\underline{\underline{F}}] = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 2 & 1 \\ 0 & 1 & 2 \end{pmatrix}$$

$$[\underline{\underline{C}}] = [\underline{\underline{F}}^T][\underline{\underline{F}}] = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 2 & 1 \\ 0 & 1 & 2 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 2 & 1 \\ 0 & 1 & 2 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 5 & 4 \\ 0 & 4 & 5 \end{pmatrix}$$

To get  $[\underline{\underline{U}}]$  we need to solve eigenvalue problem

$$\begin{vmatrix} 1-\mu & 0 & 0 \\ 0 & 5-\mu & 4 \\ 0 & 4 & 5-\mu \end{vmatrix} = (1-\mu)(5-\mu)^2 - 16(1-\mu) = 0$$

Eigenvalues:  $\mu_{1,2} = 1$     $\mu_3 = 9$

Eigen vectors:  $[\underline{u}_1] = \begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}$     $[\underline{u}_2] = \frac{1}{\sqrt{2}} \begin{bmatrix} 0 \\ -1 \\ 1 \end{bmatrix}$     $[\underline{u}_3] = \frac{1}{\sqrt{2}} \begin{bmatrix} 0 \\ 1 \\ 1 \end{bmatrix}$

$$\text{Hence: } [\underline{\underline{U}}] = \sqrt{[\underline{\underline{C}}]} = \sum_{i=1}^3 \sqrt{\mu_i} \underline{u}_i \otimes \underline{u}_i = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 2 & 1 \\ 0 & 1 & 2 \end{pmatrix}$$

$$2) \underline{U} = \sum_{i=1}^3 \lambda_i \underline{u}_i \otimes \underline{u}_i \quad \text{where}$$

$\lambda_i$ 's are principal stretches

$\underline{u}_i$ 's are right principal directions

$$\underline{C} = \underline{U}^2 = \sum_{i=1}^3 \lambda_i^2 \underline{u}_i \otimes \underline{u}_i$$

$\mu_i = \lambda_i^2$  eig. values of  $\underline{C}$  are squares of principal stretches

eigenvectors are right principal dir.

$$3) C_{KL} = F_{iK} F_{iL} \quad \text{"material strain tensor"}$$

spatial indices are contracted

### Other strain tensors

$$I) \underline{E} = \frac{1}{2}(\underline{C} - \underline{I}): \quad \text{Green-Lagrange tensor}$$

$$E_{KL} = \frac{1}{2}(C_{KL} - \delta_{KL}) \quad \text{material tensor} \Rightarrow \text{linear theory}$$

$$II) \underline{b} = \underline{F} \underline{F}^T = \underline{V}^2: \quad \text{left Cauchy-Green tensor}$$

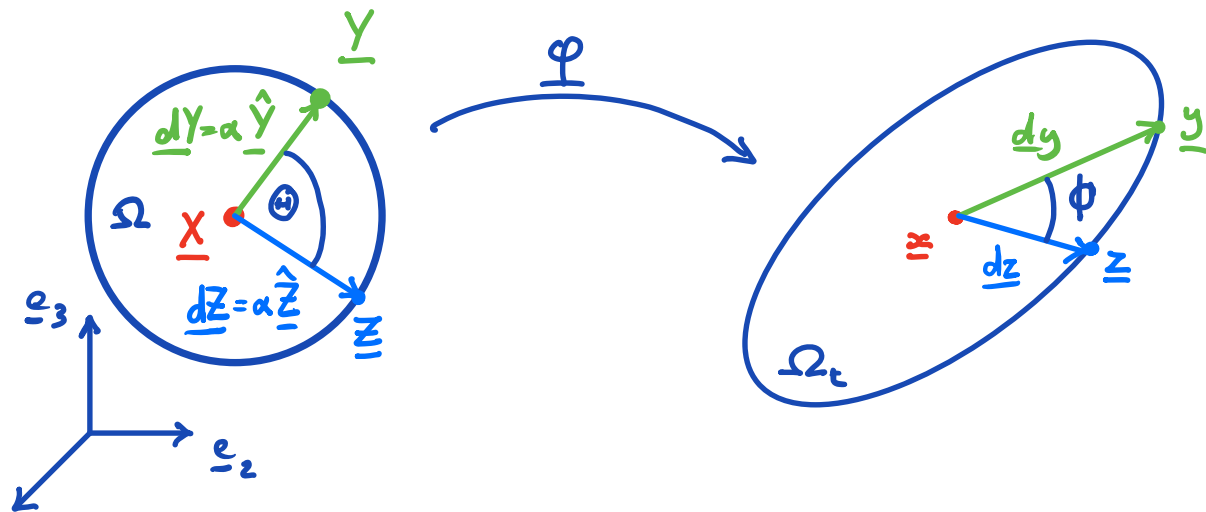
$$b_{kl} = F_{kI} F_{lI} \quad \text{"spatial tensor" (aka. Finger tensor)}$$

$$III) \underline{e} = \frac{1}{2}(\underline{I} - \underline{F}^{-T} \underline{F}^{-1}): \quad \text{Euler-Almansi tensor}$$

$$e_{kl} = \frac{1}{2}(\delta_{kl} - F_{Ik}^{-1} F_{Il}^{-1}) \quad \text{"spatial tensor"}$$

# Interpretation of $\underline{\underline{C}}$

How are changes in relative position and orientation of material points quantified by  $\underline{\underline{C}}$ ?



Consider spherical domain  $\Omega$  with radius  $\alpha > 0$  around  $\underline{X}$ . Given two unit vectors  $\hat{\underline{Y}}$  and  $\hat{\underline{Z}}$  consider the points  $\underline{Y} = \underline{X} + \alpha \hat{\underline{Y}} = \underline{X} + \underline{dY}$  and  $\underline{Z} = \underline{X} + \alpha \hat{\underline{Z}} = \underline{X} + \underline{dZ}$ .

Let  $\underline{x}, \underline{y}$  and  $\underline{z}$  denote the corresponding points  $\Omega'$  with  $\phi \in [0, \pi]$  the angle between the vectors  $\underline{dy} = \underline{y} - \underline{x}$  and  $\underline{dz} = \underline{z} - \underline{x}$ .

## Cauchy - Green strain relations

For any point  $\underline{x} \in B$  and unit vectors  $\hat{\underline{y}}$  and  $\hat{\underline{z}}$  we define  $\lambda(\hat{\underline{y}}) > 0$  and  $\theta(\hat{\underline{y}}, \hat{\underline{z}}) \in [0, \pi]$  by

$$\lambda(\hat{\underline{y}}) = \sqrt{\hat{\underline{y}} \cdot \underline{C} \hat{\underline{y}}} \quad \text{and} \quad \cos \theta(\hat{\underline{y}}, \hat{\underline{z}}) = \frac{\hat{\underline{y}} \cdot \underline{C} \hat{\underline{z}}}{\sqrt{\hat{\underline{y}} \cdot \underline{C} \hat{\underline{y}}} \sqrt{\hat{\underline{z}} \cdot \underline{C} \hat{\underline{z}}}}$$

### I. Stretches

In the limit as  $\alpha \rightarrow 0$  we have

$$\frac{|\underline{y} - \underline{x}|}{|\underline{y} - \underline{x}|} = \frac{|d\underline{y}|}{|d\underline{Y}|} \rightarrow \lambda(\hat{\underline{y}}) \quad \text{and} \quad \frac{|\underline{z} - \underline{x}|}{|\underline{z} - \underline{x}|} = \frac{|d\underline{z}|}{|d\underline{Z}|} \rightarrow \lambda(\hat{\underline{z}})$$

Therefore  $\lambda(\hat{\underline{y}})$  is the stretch in direction  $\hat{\underline{y}}$  at  $\underline{x}$ .

A stretch is the ratio of deformed to initial length.

To determine the stretch we use  $d\underline{y} = \underline{F}(d\underline{Y})$ .

$$\begin{aligned} |d\underline{y}|^2 &= d\underline{y} \cdot d\underline{y} = \underline{F} d\underline{Y} \cdot (\underline{F} d\underline{Y}) = d\underline{Y} \cdot \underline{F}^T \underline{F} d\underline{Y} = d\underline{Y} \cdot \underline{C} d\underline{Y} \\ &= \alpha^2 \hat{\underline{y}} \cdot \underline{C} \hat{\underline{y}} \end{aligned}$$

$$|d\underline{Y}|^2 = \alpha^2 \quad \text{by definition}$$

So that  $\frac{|d\underline{y}|^2}{|d\underline{Y}|^2} = \underline{\hat{y}} \cdot \underline{\underline{C}} \underline{\hat{y}} = \lambda^2(\underline{\hat{y}})$

taking square root:  $\lambda(\underline{e}) = \sqrt{\underline{\hat{y}} \cdot \underline{\underline{C}} \underline{\hat{y}}}$  ✓

If  $\underline{u}_i$  is a right-principal stretch, so that

$$(\underline{\underline{C}} - \lambda_i^2 \underline{\underline{I}}) \underline{\hat{u}}_i = 0 \quad (\text{no sum})$$

$$\underline{\hat{u}}_i \cdot \underline{\underline{C}} \underline{\hat{u}}_i - \lambda_i^2 \underline{\hat{u}}_i \cdot \underline{\hat{u}}_i = 0 \quad \underline{\hat{u}}_i \cdot \underline{\underline{C}} \underline{\hat{u}}_i = \lambda_i^2$$

note:  $\underline{\hat{u}}_i$  is the eigenvector of  $\underline{\underline{C}}$  it is capitalized because it is a material vector

then  $\lambda(\underline{\hat{u}}_i) = \lambda_i$  which justifies referring to  $\lambda_i$ 's as principal stretches.

Arguments similar to determination of principal stresses show that  $\lambda(\underline{\hat{y}})$  has extremum if  $\underline{\hat{y}} = \underline{\hat{u}}_i$ .

## II. Shear

The shear  $\gamma(\hat{\underline{Y}}, \hat{\underline{Z}})$  at  $\underline{X}$  is the change in angle between the two directions  $\hat{\underline{Y}}$  and  $\hat{\underline{Z}}$

$$\gamma(\hat{\underline{Y}}, \hat{\underline{Z}}) = \Theta(\hat{\underline{Y}}, \hat{\underline{Z}}) - \theta(\hat{\underline{Y}}, \hat{\underline{Z}})$$

where  $\Theta(\underline{e}, \underline{d})$  is the angle between  $\underline{Y}$  and  $\underline{Z}$  in the reference configuration and  $\theta(\hat{\underline{Y}}, \hat{\underline{Z}})$  is the angle between the deformed line segments  $\underline{y}$  and  $\underline{z}$  in the limit  $\alpha \rightarrow 0$  so that

$$\cos \phi \rightarrow \cos \theta(\hat{\underline{Y}}, \hat{\underline{Z}})$$

To see this consider  $\cos \phi = \frac{d\underline{y} \cdot d\underline{z}}{|d\underline{y}| |d\underline{z}|}$

where  $d\underline{y} \cdot d\underline{z} = (\underline{F} d\underline{Y}) \cdot (\underline{F} d\underline{Z})$

$$= d\underline{Y} \cdot \underline{F}^T \underline{F} d\underline{Z} = d\underline{Y} \cdot \underline{C} d\underline{Z}$$

$$= \alpha^2 \hat{\underline{Y}} \cdot \underline{C} \hat{\underline{Z}}$$

with  $|d\underline{y}| = \alpha \sqrt{\hat{\underline{Y}} \cdot \underline{C} \hat{\underline{Y}}}$  and  $|d\underline{z}| = \alpha \sqrt{\hat{\underline{Z}} \cdot \underline{C} \hat{\underline{Z}}}$

$$\text{so that } \cos \phi = \frac{d\hat{Y} \cdot \underline{\underline{C}} d\hat{Z}}{\sqrt{d\hat{Y} \cdot \underline{\underline{C}} d\hat{Y}} \sqrt{d\hat{Z} \cdot \underline{\underline{C}} d\hat{Z}}} \xrightarrow{\alpha \rightarrow 0} \cos \theta(d\hat{Y}, d\hat{Z})$$

## Components of $\underline{\underline{C}}$

Let  $C_{IJ}$  be the components of  $\underline{\underline{C}}$  in an arbitrary frame  $\{\underline{e}_I\}$ , then for any point  $\underline{x} \in B$  we have that

$$\begin{aligned} C_{II} &= \lambda^2(\underline{e}_I) \\ C_{IJ} &= \lambda(\underline{e}_I) \lambda(\underline{e}_J) \sin \gamma(\underline{e}_I, \underline{e}_J) \quad (\text{no sum}) \end{aligned}$$

$\Rightarrow$  The diagonal components of  $C$  are the squares of the stretches in coord. directions. Off diagonal components are related to shears between coordinate directions.



The expression for the diagonal components follows directly from the first Cauchy-Green strain relation

$$\lambda(\underline{Y}) = \sqrt{\underline{Y} \cdot \underline{\underline{C}} \underline{Y}} \quad \text{and} \quad C_{II} = \underline{e}_I \cdot \underline{\underline{C}} \underline{e}_I \quad (\text{no sum})$$

so that  $C_{II} = \lambda^2(\underline{e}_I)$ . ✓

For the off-diagonal components  $C_{IJ} (I \neq J)$  we start with the second Cauchy-Green strain relation

$$\cos \theta(\underline{e}_I, \underline{e}_J) = \frac{\underline{e}_I \cdot \underline{\underline{C}} \underline{e}_J}{\sqrt{\underline{e}_I \cdot \underline{\underline{C}} \underline{e}_I} \sqrt{\underline{e}_J \cdot \underline{\underline{C}} \underline{e}_J}} \quad \text{and} \quad C_{IJ} = \underline{e}_I \cdot \underline{\underline{C}} \underline{e}_J$$

so that

$$C_{IJ} = \lambda(\underline{e}_I) \lambda(\underline{e}_J) \cos \theta(\underline{e}_I, \underline{e}_J) .$$

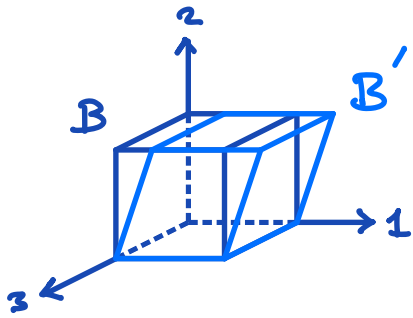
The shear between two basis vectors is

$$\gamma(\underline{e}_I, \underline{e}_J) = \underbrace{\theta(\underline{e}_I, \underline{e}_J)}_{\frac{\pi}{2}} - \theta(\underline{e}_I, \underline{e}_J)$$

so that  $C_{IJ} = \lambda(\underline{e}_I) \lambda(\underline{e}_J) \cos\left(\frac{\pi}{2} - \gamma(\underline{e}_I, \underline{e}_J)\right)$   
 $= \lambda(\underline{e}_I) \lambda(\underline{e}_J) \sin(\gamma(\underline{e}_I, \underline{e}_J))$  ✓

The components of  $\underline{\underline{C}}$  directly quantify stretch and shear unlike the components of  $\underline{\underline{E}}$ .

## Example: Simple shear



$$B = \{ \underline{X} \in \mathbb{E}^3 \mid 0 < X_i < 1 \}$$

$$\underline{x} = \underline{\varphi}(\underline{X}) = \begin{bmatrix} X_1 + \alpha X_2 \\ X_2 \\ X_3 \end{bmatrix} \quad \alpha > 0$$

"simple shear in  $\underline{e}_1$ - $\underline{e}_2$  plane"

Deformation gradient:

$$[\underline{F}] = [\nabla \underline{\varphi}] = \begin{bmatrix} \varphi_{1,1} & \varphi_{1,2} & \varphi_{1,3} \\ \varphi_{2,1} & \varphi_{2,2} & \varphi_{2,3} \\ \varphi_{3,1} & \varphi_{3,2} & \varphi_{3,3} \end{bmatrix} = \begin{bmatrix} 1 & \alpha & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$\Rightarrow$  homogeneous deformation

Cauchy-Green strain tensor:

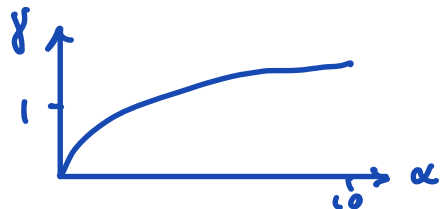
$$[\underline{C}] = \underline{F}^T \underline{F} = \begin{bmatrix} 1 & 0 & 0 \\ \alpha & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & \alpha & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} 1 & \alpha & 0 \\ \alpha & 1 + \alpha^2 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

Find the shear  $\gamma$  for direction pair  $(\underline{e}_1, \underline{e}_2)$

$$\gamma(\underline{e}_1, \underline{e}_2) = \Theta(\underline{e}_1, \underline{e}_2) - \theta(\underline{e}_1, \underline{e}_2) = \frac{\pi}{2} - \theta(\underline{e}_1, \underline{e}_2)$$

$$\cos \theta(\underline{e}_1, \underline{e}_2) = \frac{[\underline{e}_1]^T [\underline{C}] [\underline{e}_2]}{\sqrt{[\underline{e}_1]^T [\underline{C}] [\underline{e}_1]} \sqrt{[\underline{e}_2]^T [\underline{C}] [\underline{e}_2]}} = \frac{\alpha}{\sqrt{1} \sqrt{1 + \alpha^2}}$$

$$\Rightarrow \underline{\gamma}(\underline{e}_1, \underline{e}_2) = \frac{\pi}{2} - \alpha \cos\left(\frac{\alpha}{\sqrt{1 + \alpha^2}}\right)$$



Find  $\gamma(\underline{e}_1, \underline{e}_3)$  again  $\Theta(\underline{e}_1, \underline{e}_3) = \frac{\pi}{2}$

$$\cos \Theta(\underline{e}_1, \underline{e}_3) = \frac{c_{13}}{\sqrt{c_{11}} \sqrt{c_{33}}} = \frac{0}{1 \cdot 1} = 0$$

$$\gamma(\underline{e}_1, \underline{e}_3) = \frac{\pi}{2} - \underbrace{\arccos 0}_{\frac{\pi}{2}} = \underline{\underline{0}}$$

What are the extreme values of the stretch and their directions?  $\Rightarrow$  eigenvalues & vectors

$$\begin{vmatrix} 1 - \lambda^2 & \alpha & 0 \\ \alpha & 1 + \alpha^2 - \lambda^2 & 0 \\ 0 & 0 & 1 - \lambda^2 \end{vmatrix} = 0 \quad \begin{aligned} \lambda_1^2 &= 1 + \frac{\alpha^2}{2} + \alpha \sqrt{1 + \alpha^2/4} > 1 \\ \lambda_2^2 &= 1 \\ \lambda_3^2 &= 1 + \frac{\alpha^2}{2} - \alpha \sqrt{1 + \alpha^2/4} < 1 \end{aligned}$$

Principal directions:

$$[\underline{v}_1] = [\sqrt{1 + \alpha^2/4} - \alpha/2, 1, 0]$$

$$[\underline{v}_2] = [0, 0, 1]$$

(not normalized)

$$[\underline{v}_3] = [\sqrt{1 + \alpha^2/4} + \alpha/2, -1, 0]$$

$\Rightarrow \lambda_1$  is max stretch in dir  $\underline{v}_1$

$\lambda_3$  is min stretch in dir  $\underline{v}_3$

there is no stretch in dir  $\underline{e}_3$