

Lecture 15: Motion & Material Time Derivative

Logistics: - PS 6 due today

Last time: - Infinitesimal strain tensor

$$\underline{u} = \varphi(\underline{x}) - \underline{x} \quad \underline{\nabla u} = \underline{H} = \underline{F} - \underline{I}$$

$$\underline{\underline{\epsilon}} = \text{sym}(\underline{\nabla u}) = \frac{1}{2} (\underline{\nabla u} + \underline{\nabla u}^T)$$

$$\underline{\underline{C}} = \underline{\underline{I}} + 2 \underline{\underline{\epsilon}} + \underbrace{\underline{\nabla u}^T \underline{\nabla u}}_{\text{non-linear}}$$

$$\lim_{|\underline{\nabla u}| \rightarrow 0} \underline{\underline{E}} = \frac{1}{2} (\underline{\underline{C}} - \underline{\underline{I}}) = \underline{\underline{\epsilon}}$$

$$- \underline{\epsilon}_{ij} \approx \underbrace{\lambda(\underline{e}_i) - 1}_{\text{rel. change in length}} \quad \underline{\epsilon}_{ij} \approx \frac{1}{2} \underbrace{\sin \gamma(\underline{e}_i, \underline{e}_j)}_{\text{shear}} \approx \frac{1}{2} \gamma(\underline{e}_i, \underline{e}_j)$$

- Rotation - Stretch Decomposition

$$\text{finite deformation: } \underline{\underline{F}} = \underline{\underline{R}} \underline{\underline{U}}$$

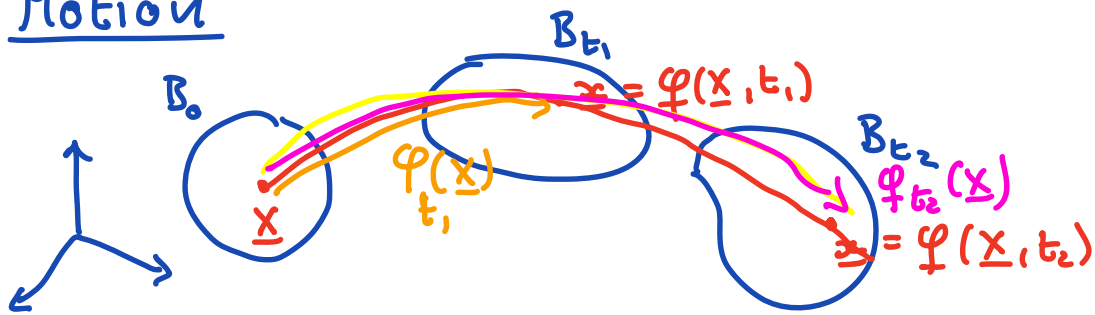
$$\text{infinitesimal defor.: } \underline{\underline{F}} \approx \underline{\underline{I}} + \underline{\underline{\epsilon}} + \underline{\underline{\omega}}$$

$$\underline{\underline{\omega}} = \text{shew}(\underline{\nabla u})$$

Today: Static \rightarrow Dynamic deformation

\Rightarrow Time derivatives

Motion



Motion is a continuous deformation of a body.
If motion is admissible \rightarrow inverse map

$$\underline{x} = \underline{\psi}(\underline{x}, t) = \underline{\varphi}^{-1}(\underline{x}, t)$$

Material field: naturally expressed in terms of \underline{x}

$$\Omega = \Omega(\underline{x}, t)$$

Spatial fields: naturally expressed in terms of \underline{x}

$$\Gamma = \Gamma(\underline{x}, t)$$

To any material field $\Omega(\underline{x}, t)$ we associate

a spatial field $\Omega_s(\underline{x}, t) = \Omega(\underline{\psi}(\underline{x}, t), t)$

we call Ω_s the spatial description of Ω

To any spatial field $\Gamma(\underline{x}, t)$ we associate
a material description: $\Gamma_m(\underline{X}, t) = \Gamma(\varphi(\underline{X}, t), t)$

Note: $\nabla_{\underline{x}}$ \rightarrow derivative with respect to material
coordinates

$\nabla_{\underline{x}}$ \rightarrow derivative with respect to spatial
coordinates

Velocity and acceleration fields

Naturally associated with particles \Rightarrow material fields

Particle initially at \underline{X} moving with motion

$$\underline{x} = \varphi(\underline{X}, t)$$

velocity:

$$\underline{V}(\underline{X}, t) = \frac{\partial}{\partial t} \varphi(\underline{X}, t) = \left. \frac{\partial \underline{x}}{\partial t} \right|_{\underline{X}}$$

acceleration:

$$\underline{A}(\underline{X}, t) = \frac{\partial^2}{\partial t^2} \varphi(\underline{X}, t) = \left. \frac{\partial^2 \underline{x}}{\partial t^2} \right|_{\underline{X}}$$

spatial descriptions are

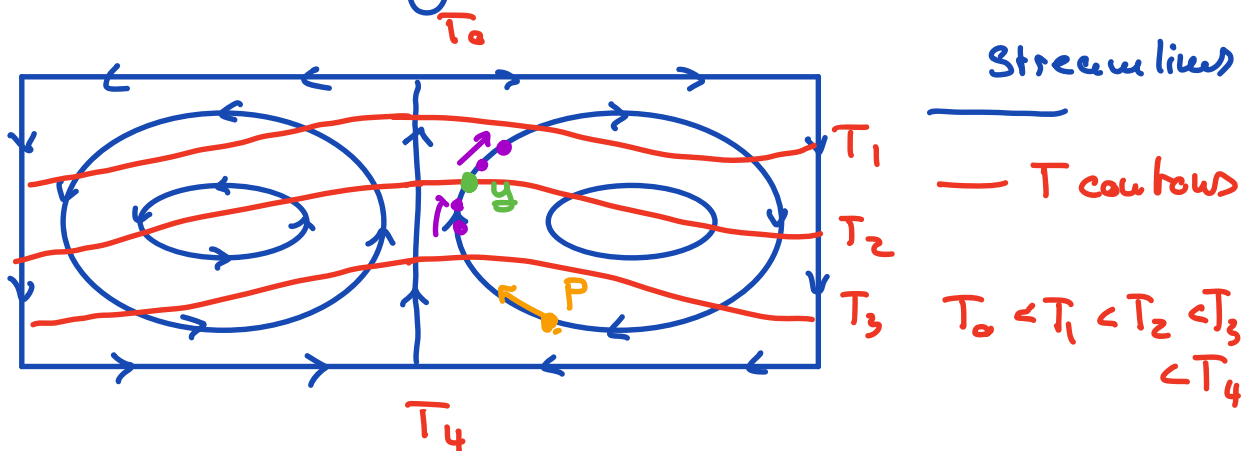
$$\underline{v}(\underline{x}, t) = \underline{V}_s(\underline{x}, t) = \frac{\partial}{\partial t} \varphi(\Psi(\underline{x}, t), t)$$

$$\underline{a}(\underline{x}, t) = \underline{A}_s(\underline{x}, t) = \frac{\partial^2}{\partial t^2} \varphi(\Psi(\underline{x}, t), t)$$

These spatial fields correspond to the velocity and acceleration of the particles whose current coordinates are \underline{x} and t .

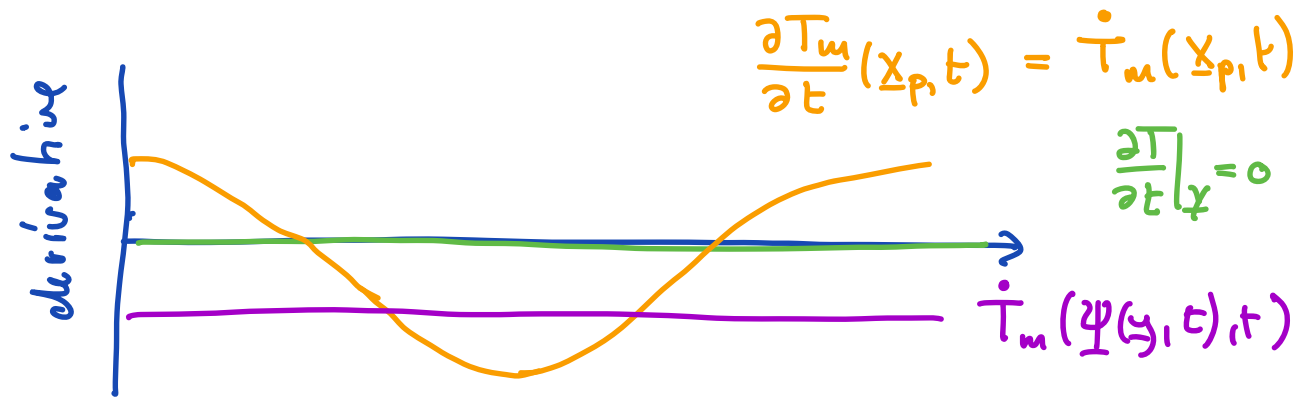
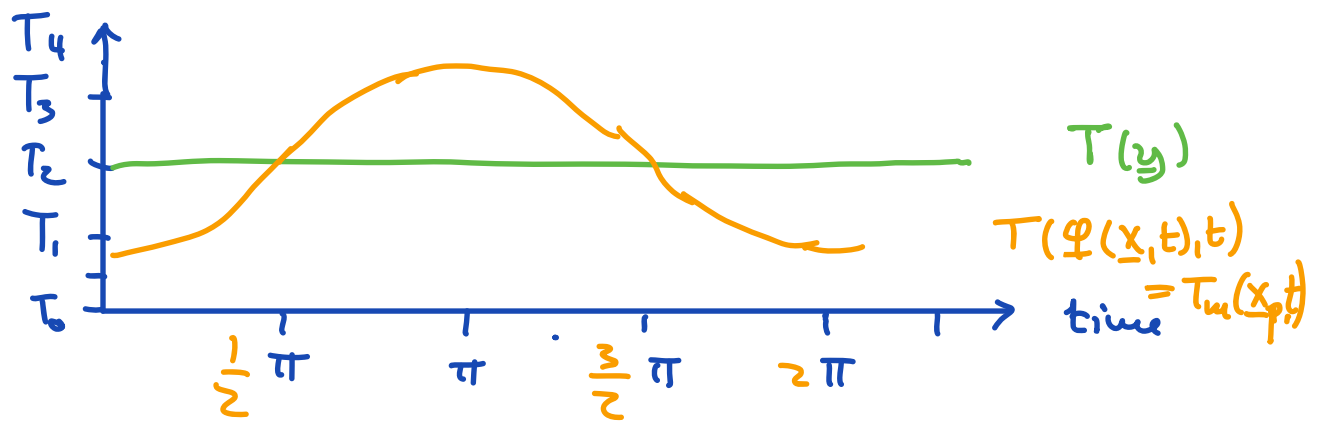
Note: Below we show that $\underline{a} \neq \frac{\partial \underline{v}}{\partial t}$

Example: Steady Convection



At steady state: $T(\underline{x}, t) = T(\underline{x})$ spatial field

Particle P. with initial location \underline{x}_p



⇒ three types of time derivatives

Different time derivatives

I) Material time derivative of a material field Ω

$$\dot{\Omega}(\underline{x}, t) = \frac{D\Omega}{Dt}(\underline{x}, t) = \frac{\partial \Omega}{\partial t} \Big|_{\underline{x}} \quad (\text{orange})$$

total, substantial, convective derivative

$\dot{\Omega}$ represents the rate of change of Ω as seen by an observer following the pathline

of a particle.

II) Spatial time derivative of spatial field Γ
Derivative with respect to t holding \underline{x} fixed

$$\frac{\partial \Gamma}{\partial t}(\underline{x}, t) = \left. \frac{\partial \Gamma}{\partial t} \right|_{\underline{x}} \quad (\text{green})$$

local time derivative

$\frac{\partial \Gamma}{\partial t}$ represents the rate of change of Γ as seen by an observer at \underline{x}

III) Material time derivative of a spatial field
Derivative of spatial field Γ with respect to time, hold \underline{x} fixed. $\underline{x} = \varphi(\underline{x}, t)$

$$\dot{\Gamma}(\underline{x}, t) = \frac{D\Gamma}{Dt}(\underline{x}, t) = \frac{\partial}{\partial t} \Gamma(\underbrace{\varphi(\underline{x}, t)}_{\underline{x}}, t) \Big|_{\underline{x} = \varphi(\underline{x}, t)}$$

$\Gamma_m(t)$

\Rightarrow two time dependencies

~

$$\frac{\partial}{\partial t} \Gamma_{\underline{x}}(\underline{x}, t) = \frac{\partial}{\partial t} \Gamma(\overline{\varphi(\underline{x}, t)}, t) \quad \text{chain rule}$$

$$= \frac{\partial \Gamma}{\partial t}(\underline{x}, t) \Big|_{\underline{x}=\varphi(\underline{x}, t)} + \frac{\partial \Gamma}{\partial x_i}(\underline{x}, t) \Big|_{\underline{x}=\varphi(\underline{x}, t)} \frac{\partial \varphi_i}{\partial t}(\underline{x}, t)$$

recognize spatial velocity field: $v_i(\underline{x}, t) \Big|_{\underline{x}=\varphi(\underline{x}, t)} = \frac{\partial \varphi_i}{\partial t}(\underline{x}, t)$

substitute

$$\frac{\partial \Gamma}{\partial t}(\overline{\varphi(\underline{x}, t)}, t) = \left[\frac{\partial \Gamma}{\partial t}(\underline{x}, t) + \frac{\partial \Gamma}{\partial x_i}(\underline{x}, t) v_i(\underline{x}, t) \right] \Big|_{\underline{x}=\varphi(\underline{x}, t)}$$

\uparrow
 material coordinates $\nabla \Gamma \cdot \underline{v}$

Expressing result in spatial coordinates

$$\dot{\Gamma}(\underline{x}, t) = \frac{\partial \Gamma}{\partial t} + \nabla_{\underline{x}} \Gamma \cdot \underline{v}$$

This is important because it allows the computation of material derivative without knowledge of the motion φ , if \underline{v} is known

For a vector $\underline{\omega}$:

$$\dot{\underline{\omega}} = \frac{\partial \underline{\omega}}{\partial t} + (\nabla_{\underline{x}} \underline{\omega}) \underline{v}$$

In many books: $(\nabla_{\underline{x}} \underline{\omega}) \underline{v} = (\underline{v} \cdot \nabla) \underline{\omega}$ (HW5)

$$\dot{T} = \frac{\partial T}{\partial t} + \underline{v} \cdot \nabla T \quad \text{scalar}$$

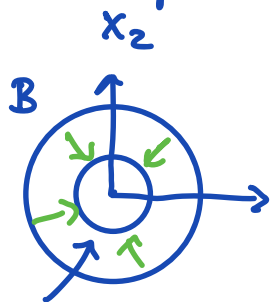
$$\dot{\underline{v}} = \frac{\partial \underline{v}}{\partial t} + (\underline{v} \cdot \nabla) \underline{v} \quad \text{vector}$$

↑ ambiguous

The spatial acceleration field is

$$\underline{a} = \dot{\underline{v}} = \frac{\partial \underline{v}}{\partial t} + (\nabla_{\underline{x}} \underline{v}) \underline{v}$$

Example: Exponential expansion/collapse



$$\underline{x} = \varphi(\underline{X}, t) = e^{-\lambda t} \underline{X} \quad \lambda > 0$$

$$\underline{X} = \psi(\underline{x}, t) = e^{\lambda t} \underline{x}$$

Be Material fields:

$$\underline{V}(\underline{x}, t) = \frac{\partial}{\partial t} \varphi(\underline{x}, t) = -\lambda \underline{x} e^{-\lambda t}$$

$$\underline{A}(\underline{x}, t) = \frac{\partial}{\partial t} \underline{V}(\underline{x}, t) = \lambda^2 \underline{x} e^{-\lambda t}$$

Spatial fields:

$$\underline{v}(\underline{x}, t) = \underline{V}_s(\underline{x}, t) = \underline{V}(\psi(\underline{x}, t), t) = -\lambda e^{\lambda t} \underline{x} e^{-\lambda t}$$

$$= -\lambda \underline{x}$$

$$\underline{a}(\underline{x}, t) = \underline{A}(\psi(\underline{x}, t), t) = \lambda^2 \underline{x}$$

Temperature field: $T_m(\underline{x}, t) = \alpha t \|\underline{x}\|$

Material time derivative:

$$\dot{T}_m = \frac{\partial}{\partial t} T_m = \alpha \|\underline{x}\|$$

calculated directly from material description

$$\text{Spatial } T: T(\underline{x}, t) = T_m(\psi(\underline{x}, t), t) = \alpha t \|\underline{x}\| e^{\lambda t}$$

Suppose we only have spatial field

$$T(\underline{x}, t) \quad \text{and} \quad \underline{v}(\underline{x}, t)$$

$$\dot{T}(\underline{x}, t) = \left(\frac{\partial T}{\partial t} \right) + \nabla_x T \cdot \underline{v}$$

$$\frac{\partial T}{\partial t} = \alpha \|\underline{x}\| e^{\lambda t} + \alpha \lambda t \|\underline{x}\| e^{\lambda t}$$

$$\nabla_x T = \alpha t e^{\lambda t} \nabla_x (\underline{x} \cdot \underline{x})^{1/2}$$

$$\nabla_x (\underline{x} \cdot \underline{x}) = 2 \underline{x}$$

$$= \alpha t e^{\lambda t} \frac{\underline{x}}{\|\underline{x}\|}$$

Put it all together

$$\dot{T} = \underbrace{\alpha \|\underline{x}\| e^{\lambda t} + \alpha \lambda t \|\underline{x}\| e^{\lambda t}}_{\frac{\partial T}{\partial t}} + \underbrace{\alpha t e^{\lambda t} \frac{\underline{x}}{\|\underline{x}\|}}_{\nabla_x T} \cdot (-\lambda \underline{x})$$

$$= \alpha \|\underline{x}\| e^{\lambda t} + \cancel{\alpha \lambda t \|\underline{x}\| e^{\lambda t}} - \alpha \lambda t e^{\lambda t} \underbrace{\frac{\underline{x} \cdot \underline{x}}{\|\underline{x}\|}}_{\|\underline{x}\|}$$

$$\dot{T}(\underline{x}, t) = \underline{\alpha \|\underline{x}\| e^{\lambda t}}$$

from above $\dot{T}_u(\underline{X}, t) = \alpha \|\underline{X}\|$

$$\dot{T} = \dot{T}_u(\psi(\underline{x}, t)) = \alpha \|e^{\lambda t} \underline{x}\| = \underline{\alpha \|\underline{x}\| e^{\lambda t}}$$