

Lecture 22: Objectivity and Material constraints

Logistics: HW8 due

HW9 will be posted

Projects due Dec 6th

Last time: • Constitutive laws

$$\underline{\underline{G}}(\underline{\underline{A}}) = \underline{\underline{C}}\underline{\underline{A}} = \lambda \text{tr}(\underline{\underline{A}})\underline{\underline{I}} + 2\mu \text{sym}(\underline{\underline{A}})$$

$$A \rightarrow \nabla_{\underline{v}} \text{ or } \nabla_{\underline{u}}$$

• What is a fourth-order tensor?

4-th order dyadic product

$$(\underline{a} \otimes \underline{b} \otimes \underline{c} \otimes \underline{d})_{\underline{\underline{T}}} = (\underline{c} \cdot \underline{\underline{T}} \underline{d}) \underline{a} \otimes \underline{b}$$

4-th order tensor

$$\underline{\underline{C}} = C_{ijkl} \underline{e}_i \otimes \underline{e}_j \otimes \underline{e}_k \otimes \underline{e}_l$$

$$C_{ijkl} = \underline{e}_i \cdot \underline{\underline{C}} (\underline{e}_k \otimes \underline{e}_l) \underline{e}_j$$

$$\underline{\underline{U}} = \underline{\underline{C}} \underline{\underline{T}} \Rightarrow U_{ij} = C_{ijkl} T_{kl}$$

• Major & minor symmetries

Today: Objectivity, Representation Theorem, Constraints

Change of observer

Lecture 6 we discussed change in basis

$$\{\underline{e}_i\} \text{ and } \{\underline{e}'_i\} \quad \underline{v} = \underline{Q} \underline{v}' \quad \underline{s} = \underline{Q} \underline{s}' \underline{Q}^T$$

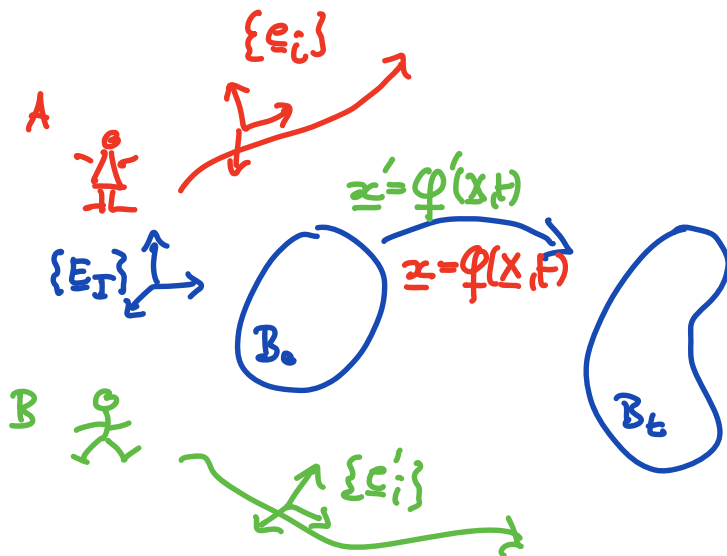
whr \underline{Q} is a rotation (change in basis tensor)

1) orthonormal: $\underline{Q}^T \underline{Q} = \underline{I} = \underline{Q} \underline{Q}^T \Rightarrow \underline{Q}^T = \underline{Q}^{-1}$

2) $\det(\underline{Q}) = 1$

Change in basis is passive change in frame.

Change in observer (active change in frame)



Material frame
is common but
spatial frame is
different.

(Note: assume same clock)

Change in observer cannot induce a deformation $\Rightarrow \varphi$ and φ' must be related by a superposed rigid motion.

$$\underline{x}' = \underline{Q}(t) \underline{x} + \underline{c}(t) \quad \underline{x} = \varphi(\underline{X}, t)$$

$$\underline{x}' = \underline{Q}(t) \varphi(\underline{X}, t) + \underline{c}(t) \quad \text{Eulerian transformation.}$$

\underline{Q} = rotation \underline{c} = translation.

Objective description of forces and deformations cannot depend on the observer.

Effect of superposed rigid motion of kinematic quantities

$$\underline{x} = \varphi(\underline{X}, t) \quad \nabla_{\underline{x}} \varphi = \underline{F}$$

$$\underline{x}' = \varphi'(\underline{X}, t) = \underline{Q} \varphi(\underline{X}, t) + \underline{c}(t)$$

$$\nabla \varphi' = \underline{Q} \nabla \varphi = \underline{Q} \underline{F} = \underline{F}' \quad \underline{F} = \underline{R} \underline{U}$$

Right Cauchy-Green strain tensor: $\underline{C} = \underline{F}^T \underline{F}$

$$\underline{C}' = \underline{F}'^T \underline{F}' = (\underline{Q} \underline{F})^T (\underline{Q} \underline{F}) = \underline{F}^T \underbrace{\underline{Q}^T \underline{Q}}_{\underline{I}} \underline{F} = \underline{F}^T \underline{F} = \underline{C}$$

$\Rightarrow \underline{\underline{C}}$ is not affected by change in ⁼ observer \Rightarrow objective
 material tensor C_{IJ} naturally objective

What about spatial tensors?

Axiom of frame indifference

Field ϕ , $\underline{\omega}$ and $\underline{\underline{S}}$ are called frame indifferent or objective if for all superposed rigid motions $\underline{x}' = \underline{Q} \underline{x} + \underline{c}$ we have

$\phi'(\underline{x}', t) = \phi(\underline{x}, t)$	scalar field
$\underline{\omega}'(\underline{x}', t) = \underline{Q} \underline{\omega}(\underline{x}, t)$	vector field
$\underline{\underline{S}}'(\underline{x}', t) = \underline{Q} \underline{\underline{S}}(\underline{x}, t) \underline{Q}^T$	tensor field

see lecture 6

Is spatial velocity gradient objective?

Lecture 16: $\underline{\underline{L}} = \nabla_{\underline{x}} \underline{\underline{\sigma}} = \underline{\underline{\dot{F}}} \underline{\underline{F}}^{-1}$
 $\underline{\underline{L}}' = \nabla_{\underline{x}'} \underline{\underline{\sigma}}' = \underline{\underline{\dot{F}}}' \underline{\underline{F}}'^{-1}$

$$\underline{\underline{F}}' = \underline{\underline{Q}} \underline{\underline{F}} \quad \underline{\underline{\ell}}' = \underline{\underline{Q}} \underline{\underline{\ell}} \underline{\underline{Q}}^T$$

$$\underline{\underline{\dot{F}}}' = \frac{d}{dt} (\underline{\underline{Q}} \underline{\underline{F}}) = \underline{\underline{Q}} \underline{\underline{\dot{F}}} + \underline{\underline{\dot{Q}}} \underline{\underline{F}}$$

$$\underline{\underline{F}}'^{-1} = (\underline{\underline{Q}} \underline{\underline{F}})^{-1} = \underline{\underline{F}}^{-1} \underline{\underline{Q}}^{-1} = \underline{\underline{F}}^{-1} \underline{\underline{Q}}^T$$

$$\begin{aligned} \underline{\underline{\ell}}' &= \underline{\underline{\dot{F}}}' \underline{\underline{F}}'^{-1} = (\underline{\underline{Q}} \underline{\underline{\dot{F}}} + \underline{\underline{\dot{Q}}} \underline{\underline{F}}) \underline{\underline{F}}^{-1} \underline{\underline{Q}}^T \\ &= \underline{\underline{Q}} \underbrace{\underline{\underline{\dot{F}}} \underline{\underline{F}}^{-1}}_{\underline{\underline{\ell}}} \underline{\underline{Q}}^T + \underline{\underline{\dot{Q}}} \underbrace{\underline{\underline{F}} \underline{\underline{F}}^{-1}}_{\underline{\underline{I}}} \underline{\underline{Q}}^T \end{aligned}$$

$$\underline{\underline{\ell}}' = \underline{\underline{Q}} \underline{\underline{\ell}} \underline{\underline{Q}}^T + \underline{\underline{\dot{Q}}} \underline{\underline{Q}}^T \quad \text{not objective}$$

⇒ that's why $\nabla_x \underline{\underline{\sigma}}$ is not used in constitutive laws

The non-objective term $\underline{\underline{\Omega}} = \underline{\underline{\dot{Q}}} \underline{\underline{Q}}^T$

it represents the rigid body angular velocity change between the observers.

$$\text{HW9} \rightarrow \underline{\underline{\Omega}} = -\underline{\underline{\Omega}}^T \quad \text{is skew symmetric} \\ \rightarrow \text{rotation}$$

⇒ base our constitutive laws on symmetric part

$$\text{of } \underline{\underline{\ell}} \quad \underline{\underline{d}} = \text{sym}(\underline{\underline{\ell}}) = \frac{1}{2} (\nabla \underline{\underline{v}} + \nabla \underline{\underline{v}}^T)$$

$$\text{sym}(\underline{\underline{d}}') = \underline{\underline{Q}} \text{sym}(\underline{\underline{d}}) \underline{\underline{Q}}^T + \cancel{\text{sym}(\underline{\underline{\omega}})}$$

$$\underline{\underline{d}}' = \underline{\underline{Q}} \underline{\underline{d}} \underline{\underline{Q}}^T$$

rate of strain tensor is objective

⇒ used in constitutive laws

Material frame indifferent functions

Field: $\phi(\underline{x}, t)$ scalar

$\underline{\omega}(\underline{x}, t)$

$\underline{\underline{d}}(\underline{x}, t)$

fields because they depend directly on \underline{x}

Constitutive functions are not fields

but they depend on fields as input.

internal energy: $u(\underline{x}, t) = \hat{u}(\rho(\underline{x}, t), \theta(\underline{x}, t)) = \hat{u}(\rho, \theta)$

$u \rightarrow$ field $\hat{u} \rightarrow$ is constitutive function

heat flow: $q(\underline{x}, t) = \hat{q}(\theta(\underline{x}, t)) = \hat{q}(\theta)$

Cauchy stress: $\underline{\underline{\sigma}}(\underline{x}, t) = \underline{\underline{\hat{\sigma}}}(\rho(\underline{x}, t), \theta(\underline{x}, t), \underline{\underline{d}}(\underline{x}, t))$
 $= \underline{\underline{\hat{\sigma}}}(\rho, \theta, \underline{\underline{d}})$

Constitutive functions: $\hat{u}(\rho, \theta)$, $\hat{q}(\theta)$, $\underline{\underline{\hat{\sigma}}}(\rho, \theta, \underline{\underline{d}})$

$\hat{u} = \rho c_p \theta$ $u(\underline{x}, t) = u'(\underline{x}', t)$

As such constitutive functions are not directly dependent on frame but their input fields are

Consider two frames $\{\underline{e}_i\}$ and $\{\underline{e}'_i\}$

$\underline{\underline{\sigma}}'(\underline{x}', t) = \underline{\underline{\hat{\sigma}}}(\rho', \theta', \underline{\underline{d}}')$ $= \underline{\underline{\hat{\sigma}}}(\rho', \theta', \underline{\underline{Q}} \underline{\underline{d}} \underline{\underline{Q}}^T)$

$\rho' = \rho$
 $\underline{\underline{\sigma}}' = \underline{\underline{Q}} \underline{\underline{\sigma}} \underline{\underline{Q}}^T$

$\underline{\underline{\hat{\sigma}}}(\rho', \theta', \underline{\underline{Q}} \underline{\underline{d}} \underline{\underline{Q}}^T) = \underline{\underline{Q}} \underline{\underline{\hat{\sigma}}}(\rho, \theta, \underline{\underline{d}}) \underline{\underline{Q}}^T$

condition for $\underline{\underline{\hat{\sigma}}}$ to be invariant/objective

Isotropic functions

Functions that are frame invariant are called isotropic.

functions: $\hat{\phi}$ $\hat{\omega}$ $\hat{\sigma}$

inputs: θ \underline{v} \underline{s}

Following requirements: of for isotropic functions

$$\hat{\phi}(\theta) = \hat{\phi}(\theta) \quad \hat{\phi}(\underline{Q}\underline{v}) = \hat{\phi}(\underline{v}) \quad \hat{\phi}(\underline{Q}\underline{s}\underline{Q}^T) = \hat{\phi}(\underline{s})$$

$$\hat{\omega}(\theta) = \underline{Q}\hat{\omega}(\theta) \quad \hat{\omega}(\underline{Q}\underline{v}) = \underline{Q}\hat{\omega}(\underline{v}) \quad \hat{\omega}(\underline{Q}\underline{s}\underline{Q}^T) = \underline{Q}\hat{\omega}(\underline{s})$$

$$\hat{\sigma}(\theta) = \underline{Q}\hat{\sigma}(\theta)\underline{Q}^T \quad \hat{\sigma}(\underline{Q}\underline{v}) = \underline{Q}\hat{\sigma}(\underline{v})\underline{Q}^T \quad \hat{\sigma}(\underline{Q}\underline{s}\underline{Q}^T) = \underline{Q}\hat{\sigma}(\underline{s})\underline{Q}^T$$

Examples:

$$1) \hat{\phi}(\underline{s}) = \det(\underline{s})$$

$$\hat{\phi}(\underline{Q}\underline{s}\underline{Q}^T) = \det(\underline{Q}\underline{s}\underline{Q}^T) = \det(\underline{Q})\det(\underline{s})\det(\underline{Q}^T)$$
$$\hat{\phi}(\underline{s}') = \det(\underline{s}) = \hat{\phi}(\underline{s})$$

$$\hat{\phi}(\underline{Q}\underline{s}\underline{Q}^T) = \hat{\phi}(\underline{s})$$

$$2) \hat{u}(\underline{v}, \underline{A}) = \underline{A} \underline{v}$$

$$\hat{u}(\underline{Q} \underline{v}, \underline{Q} \underline{A} \underline{Q}^T) = (\underline{Q} \underline{A} \underline{Q}^T) (\underline{Q} \underline{v}) = \underline{Q} \underline{A} \underline{Q}^T \underline{Q} \underline{v} = \underline{Q} \underline{A} \underline{v} = \underline{Q} \hat{u}(\underline{v}, \underline{A})$$

Representation of isotropic tensor functions

An isotropic function $\underline{G}(\underline{A}) : \mathcal{V}^2 \rightarrow \mathcal{V}^2$ that maps sym. tensors into sym. tensors must have the following form.

$$\underline{G}(\underline{A}) = \alpha_0(\underline{I}_A) \underline{I} + \alpha_1(\underline{I}_A) \underline{A} + \alpha_2(\underline{I}_A) \underline{A}^2$$

Rivlin-Ericksen
Repres. Thm

where $\alpha_0, \alpha_1, \alpha_2$ are functions of

the set of principal invariants of \underline{A}

$$\underline{I}_A = \{ I_1(\underline{A}), I_2(\underline{A}), I_3(\underline{A}) \}$$

- \underline{G} is sym. if \underline{A} is sym.
- To see \underline{G} is isotropic assume $\alpha_0, \alpha_1, \alpha_2 = \text{const}$

$$\begin{aligned} \underline{G}(\underline{Q} \underline{A} \underline{Q}^T) &= \alpha_0 \underline{I} + \alpha_1 \underline{Q} \underline{A} \underline{Q}^T + \alpha_2 \underline{Q} \underline{A} \underline{Q}^T \underline{Q} \underline{A} \underline{Q}^T \\ &= \alpha_0 \underline{Q} \underline{I} \underline{Q}^T + \alpha_1 \underline{Q} \underline{A} \underline{Q}^T + \alpha_2 \underline{Q} \underline{A}^2 \underline{Q}^T \end{aligned}$$

$$\begin{aligned}
 &= \underline{\underline{Q}} (\alpha_0 \underline{\underline{I}} + \alpha_1 \underline{\underline{A}} + \alpha_2 \underline{\underline{A}}^2) \underline{\underline{Q}}^T \\
 &= \underline{\underline{Q}} G(A) \underline{\underline{Q}}^T \quad \checkmark \quad \text{isotropic/objective}
 \end{aligned}$$

This is the most general form of a constitutive eqn.

Isotropic 4th order tensor

If $\underline{\underline{G}}(\underline{\underline{A}})$ is linear function

$$\underline{\underline{G}}(\underline{\underline{A}}) = \underline{\underline{C}} \underline{\underline{A}}$$

where $\underline{\underline{C}}$ is 4th order tensor

If we also require:

1) $\underline{\underline{C}} \underline{\underline{A}}$ is sym. for every sym. $\underline{\underline{A}}$.

2) $\underline{\underline{C}} \underline{\underline{W}} = \underline{\underline{0}}$ for every skew-sym. $\underline{\underline{W}}$

Then there are two scalars λ, μ such that

$$\underline{\underline{G}}(\underline{\underline{A}}) = \underline{\underline{C}} \underline{\underline{A}} = \lambda \operatorname{tr}(\underline{\underline{A}}) \underline{\underline{I}} + 2\mu \operatorname{sym}(\underline{\underline{A}})$$

This follows from representation theorem.

$$\underline{\underline{G}}(\underline{\underline{A}}) = \alpha_0(\underline{\underline{I}}_A) \underline{\underline{I}} + \alpha_1(\underline{\underline{I}}_A) \underline{\underline{A}} + \alpha_2(\underline{\underline{I}}_A) \underline{\underline{A}}^2$$

where $\underline{\underline{I}}_A = \left\{ \text{tr}(\underline{\underline{A}}), \frac{1}{2} [\text{tr}(\underline{\underline{A}})^2 - \text{tr}(\underline{\underline{A}}^2)], \det \underline{\underline{A}} \right\}$

Lecture 6

\uparrow
 $\underline{\underline{I}}_1$

\uparrow
 $\underline{\underline{I}}_2$

\uparrow
 $\underline{\underline{I}}_3$

only $\text{tr}(A)$ is a linear function

for $\underline{\underline{G}}(\underline{\underline{A}})$ to be linear in $\underline{\underline{A}}$:

$$\alpha_2 = 0 \quad \alpha_1 = \text{const.} = c_2$$

$$\alpha_0 = c_0 \text{tr}(\underline{\underline{A}}) + c_1 \quad \alpha_1 = c_2$$

where c_0, c_1, c_2 are const.

$$\underline{\underline{G}}(\underline{\underline{0}}) = \underline{\underline{0}} \Rightarrow c_1 = 0$$

$$\text{set } c_0 = \lambda \quad c_2 = 2\mu$$

$$\Rightarrow \underline{\underline{G}}(\underline{\underline{A}}) = \underline{\underline{G}} \underline{\underline{A}} = \lambda \text{tr}(\underline{\underline{A}}) \underline{\underline{I}} + 2\mu \underline{\underline{A}}$$

$$\text{if } \underline{\underline{G}}(\underline{\underline{w}}) = 0 \quad (\text{tr}(\underline{\underline{w}}) = 0)$$

$$\Rightarrow \underline{\underline{G}}(\underline{\underline{A}}) = \underline{\underline{G}} \underline{\underline{A}} = \lambda \text{tr}(\underline{\underline{A}}) \underline{\underline{I}} + 2\mu \text{sym}(\underline{\underline{A}})$$

most general linear constitutive law!

Example linear elasticity : $\underline{\underline{A}} = \nabla \underline{u}$ $\text{tr}(\nabla \underline{u}) = \nabla \cdot \underline{u}$

$$\underline{\underline{\epsilon}} = \lambda (\nabla \cdot \underline{u}) \underline{\underline{I}} + \mu (\nabla \underline{u} + \nabla \underline{u}^T) \quad \checkmark$$