



# Lecture 19: Motion and Material Derivative

- Logistics:
- HW 7 due Thu
  - HW 6 → please complete
  - HW 5 → graded

Last time: - Infinitesimal strain tensors




$\underline{\underline{\epsilon}} = \frac{1}{2} (\underline{\nabla} \underline{u} + \underline{\nabla} \underline{u}^T)$        $\underline{u} = \underline{\varphi} - \underline{x}$



$\lim_{|\underline{\nabla} \underline{u}| \rightarrow 0} \underline{\underline{E}} = \underline{\underline{\epsilon}}$        $\underline{\underline{E}} = \frac{1}{2} (\underline{\underline{C}} - \underline{\underline{I}})$

Components:  $\underline{\epsilon}_{ii} \approx \lambda(\underline{e}_i) - 1$



$\underline{\epsilon}_{ij} \approx \frac{1}{2} \sin \gamma(\underline{e}_i, \underline{e}_j)$

- Linearization:

$$\underline{\underline{C}} \approx \underline{\underline{I}} + \underline{\nabla} \underline{u} + \underline{\nabla} \underline{u}^T \quad \text{sym}(\underline{\nabla} \underline{u})$$

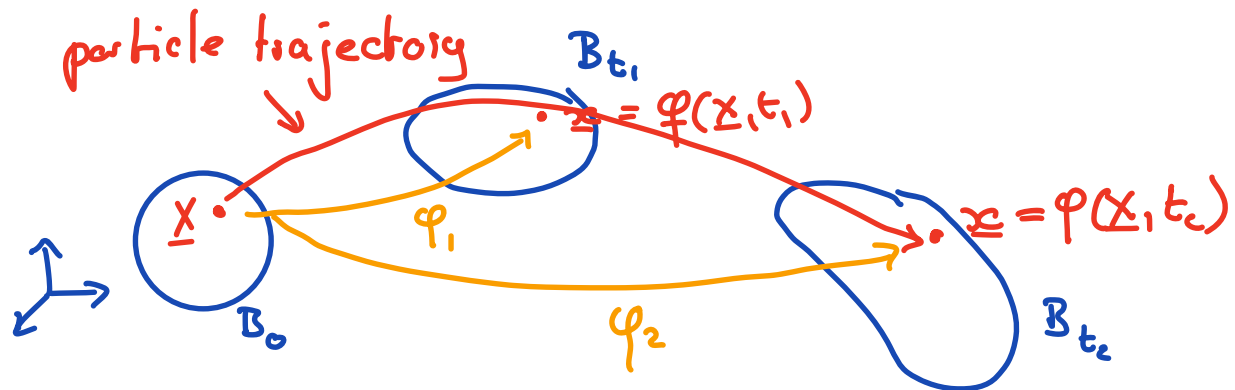
$$\underline{\underline{U}} \approx \underline{\underline{I}} + \frac{1}{2} (\underline{\nabla} \underline{u} + \underline{\nabla} \underline{u}^T) = \underline{\underline{I}} + \underline{\underline{\epsilon}}$$

$$\underline{\underline{R}} \approx \underline{\underline{I}} + \frac{1}{2} (\underline{\nabla} \underline{u} - \underline{\nabla} \underline{u}^T) = \underline{\underline{I}} - \underline{\underline{\omega}}$$

$$\underline{\underline{F}} = \underline{\underline{R}} \underline{\underline{U}} \approx \underline{\underline{I}} + \underline{\underline{\epsilon}} + \underline{\underline{\omega}} \quad \text{skew}(\underline{\nabla} \underline{u})$$

Today: Motion →  $\underline{\varphi}(\underline{x}, t)$   
 Material time derivative

# Motion



- $B_0$ :
- undeformed
  - reference
  - material
  - initial

- $B_t$ :
- deformed
  - spatial
  - current

Continuous deformation of body over time is a motion

$$\underline{x}(t) = \varphi_t(\underline{x}) = \varphi(\underline{x}, t)$$

assume  $\varphi_t$  smooth  $\Rightarrow$  admissible

$\Rightarrow$  inverse motion

$$\underline{x} = \underline{\psi}_t(\underline{x}) = \underline{\psi}(\underline{x}, t)$$

## Material and spatial fields

Temperature is naturally a spatial field:  $T(\underline{x}, t)$

Velocity is naturally a material field:  $V(\underline{X}, t)$

$\varphi_t$  &  $\psi_t$  : material  $\leftrightarrow$  spatial

Material field:  $\Omega = \Omega(\underline{X}, t)$

Spatial field:  $\Gamma = \Gamma(\underline{x}, t)$

$$\underline{X} = \psi_t(\underline{x})$$

$$\underline{x} = \varphi(\underline{X}, t)$$

Spatial description of material field

$$\Omega_s(\underline{x}, t) = \Omega(\psi(\underline{x}, t), t)$$

Material description of a spatial field

$$\Gamma_m(\underline{X}, t) = \Gamma(\varphi(\underline{X}, t), t)$$

Coordinate derivatives:

Material coordinates:  $\nabla_{\underline{X}} = \text{Grad, Div, Curl}$

Spatial coordinates:  $\nabla_{\underline{x}}$

## Velocity and Acceleration fields

The velocity and acceleration of a material particle  $\underline{x}$  at time  $t$  due to the motion  $\varphi(\underline{x}, t)$  are:

$$V(\underline{x}, t) = \frac{\partial}{\partial t} \underbrace{\varphi(\underline{x}, t)}_{\underline{x}} = \left. \frac{\partial \underline{x}}{\partial t} \right|_{\underline{x}}$$

$$A(\underline{x}, t) = \frac{\partial^2}{\partial t^2} \varphi(\underline{x}, t) = \left. \frac{\partial^2 \underline{x}}{\partial t^2} \right|_{\underline{x}}$$

naturally material fields

Spatial descriptions:

$$\underline{v}(\underline{x}, t) = V_s(\underline{x}, t) = \frac{\partial}{\partial t} \varphi(\psi(\underline{x}, t), t)$$

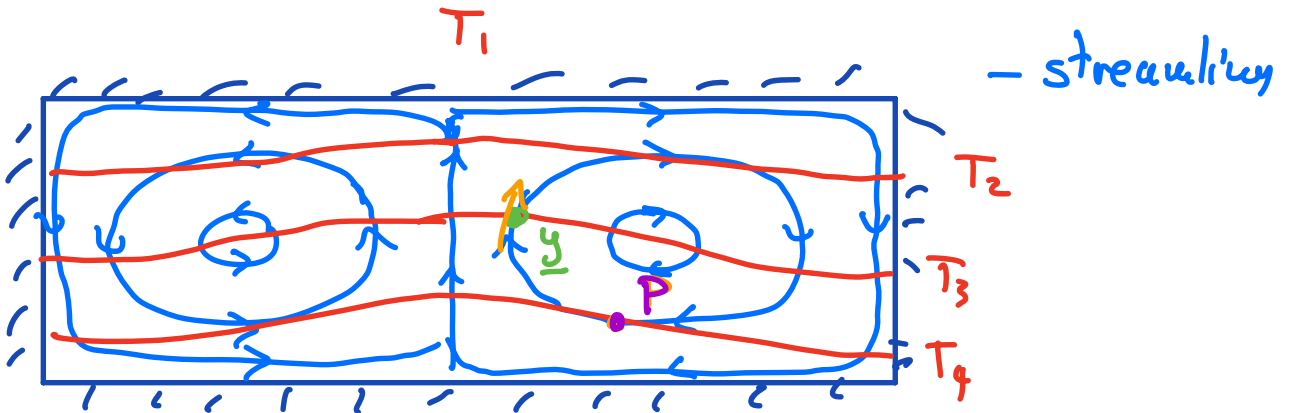
$$\underline{a}(\underline{x}, t) = A_s(\underline{x}, t) = \frac{\partial^2}{\partial t^2} \varphi(\psi(\underline{x}, t), t)$$

The spatial fields  $\underline{v}$  and  $\underline{a}$  correspond to the material particle that passes through  $\underline{x}$  at  $t$ .

Note:  $\underline{a} \neq \frac{\partial \underline{v}}{\partial t}$

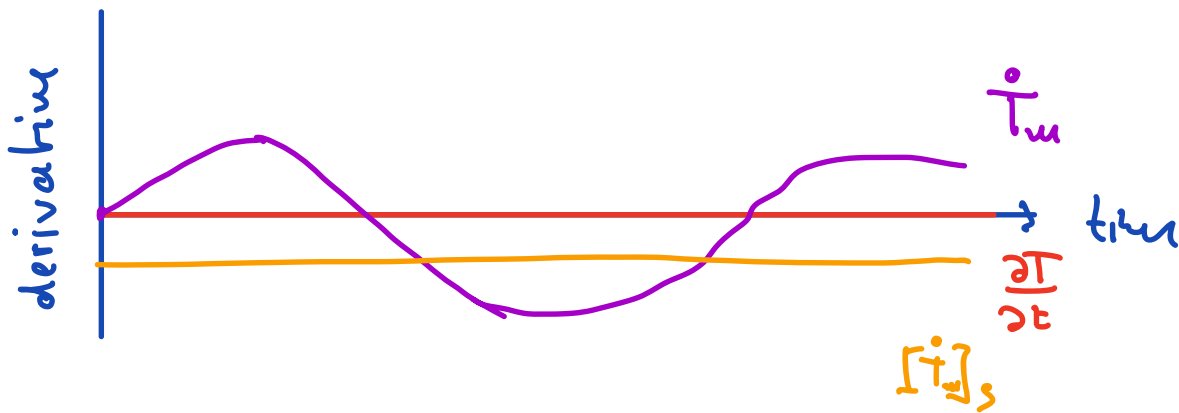
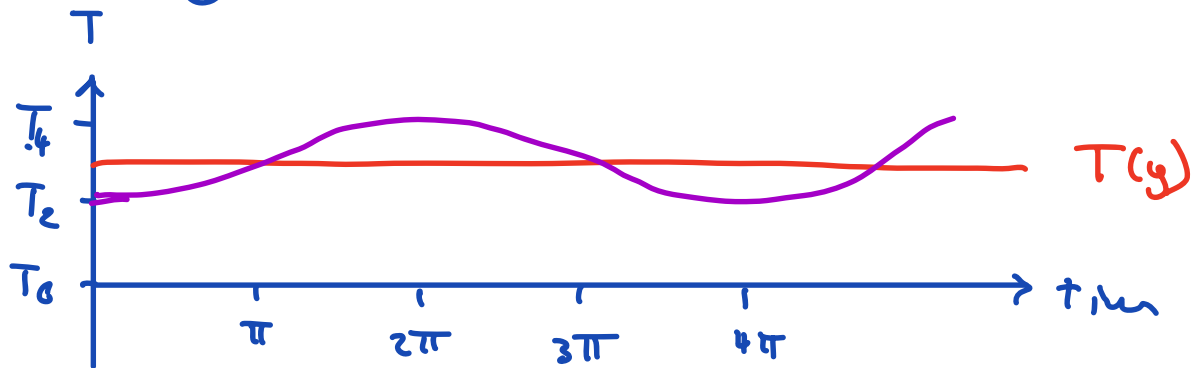
# Example: Steady Convection

$$T_4 > T_1$$



$$\dot{T}(\underline{x}, t) = \cancel{\frac{\partial T}{\partial t}} + \underline{\nabla T} \cdot \underline{v}$$

Steady state:  $T(\underline{x}, t) = T(\underline{x})$  (spatial field)



4)  $T(y) \quad \frac{\partial T}{\partial t} = 0$

2) Particle P initially at  $\underline{x}$  now at  $\underline{x} = \varphi(\underline{x}, t)$   
passes through  $y$  every  $2\pi$

$$T(\varphi(\underline{x}, t)) = T_m(\underline{x}, t) \quad \text{oscillates}$$

between  $T_2$  &  $T_4$

Its derivative "material derivative"

$$\dot{T}_m(\underline{x}, t) \neq 0$$

3) Consider particles passing through  $y$   
with initial locations  $\underline{y} = \psi(y, t)$ .

What is change in temperature these  
particles experience as they pass through

$$y? \quad \dot{T}_m(\psi(y, t)) = [\dot{T}_m]_s$$

spatial representation of material time derivative

$$\underline{Q} = \underline{x} t + \exp(-\lambda t)$$

$$\dot{\underline{Q}} = \underline{x} + \lambda \exp(-\lambda t)$$

## Different time derivatives

I, Material time derivative of a material field

$$\dot{\Omega}(\underline{x}, t) = \frac{D\Omega}{Dt}(\underline{x}, t) = \left. \frac{\partial \Omega}{\partial t} \right|_{\underline{x}}$$

derivative of  $\Omega$  with respect to  $t$  holding  $\underline{x}$  fixed

Called: total, substantial, convective material derivative

$\dot{\Omega}$  represents rate of change of  $\Omega$  seen by an observer following the particle.

II Spatial time derivative of a spatial field

$$\left. \frac{\partial \Gamma}{\partial t}(\underline{x}, t) \right|_{\underline{x}} = \frac{\partial \Gamma}{\partial t}$$

local time derivative

Rate of change of  $\Gamma$  seen by stationary observer at  $\underline{x}$

### III Material time derivative of spatial field

Derivative of scalar field  $\Gamma$  with respect to time  $t$ , holding  $\underline{x}$  fixed.

$$\Rightarrow \underline{x} = \varphi(\underline{X}, t)$$

$$\dot{\Gamma}(\underline{x}, t) = \frac{\partial}{\partial t} \underbrace{\Gamma(\varphi(\underline{X}, t), t)}_{\Gamma_m(\underline{X}, t)} \Big|_{\underline{x} = \varphi(\underline{x}, t)}$$

$\Rightarrow$  two time dependencies, one explicit other through the motion  $\varphi(\underline{x}, t)$

By chain rule:

$$\frac{\partial}{\partial t} \Gamma(\varphi(\underline{x}, t), t) \Big|_{\underline{x} = \varphi(\underline{x}, t)} = \frac{\partial \Gamma}{\partial t} \Big|_{\underline{x} = \varphi(\underline{x}, t)} + \frac{\partial \Gamma}{\partial x_i}(\underline{x}, t) \Big|_{\underline{x} = \varphi(\underline{x}, t)} \frac{\partial \varphi_i}{\partial t}(\underline{x}, t)$$

recognize spatial velocity:  $v_i(\underline{x}, t) \Big|_{\underline{x} = \varphi(\underline{x}, t)} = \frac{\partial \varphi_i}{\partial t}$

substitution

$$\dot{\Gamma}(\underline{x}, t) = \left[ \frac{\partial \Gamma}{\partial t}(\underline{x}, t) + \frac{\partial \Gamma}{\partial x_i}(\underline{x}, t) v_i(\underline{x}, t) \right] \Big|_{\underline{x} = \varphi(\underline{x}, t)}$$



Expressing this in spatial coordinates

$$\dot{\Gamma}(\underline{x}, t) = \frac{\partial \Gamma}{\partial t} + \frac{\partial \Gamma}{\partial x_i} v_i$$

Let  $\varphi(\underline{x}, t)$  be motion with spatial velocity field  $\underline{v}$  and scalar spatial field  $\phi(\underline{x}, t)$  vector spatial field  $\underline{w}$  the the spatial representations of the material time derivatives are:

$$\begin{aligned} \dot{\phi} &= \frac{\partial \phi}{\partial t} + \nabla_x \phi \cdot \underline{v} \leftarrow \text{familiar} \\ \dot{\underline{w}} &= \frac{\partial \underline{w}}{\partial t} + (\nabla_x \underline{w}) \underline{v} \end{aligned}$$

Fluid mechanics:  $(\nabla_x \underline{w}) \underline{v} = \underline{v} \cdot \nabla \underline{w}$

These results are important because they allow the computation of  $\dot{\phi}$  and  $\dot{\underline{w}}$  with out knowledge of  $\varphi$  if  $\underline{v}$  is known!  
 $\Rightarrow$  in fluid mechanics we now see  $\varphi$

The spatial acceleration can be computed as

$$\underline{a}(\underline{x}, t) = \dot{\underline{v}}(\underline{x}, t) = \frac{\partial \underline{v}}{\partial t} + (\underbrace{\nabla_{\underline{x}} \underline{v}}_{\text{non-linear}}) \underline{v}$$

⇒ basic non-linearity in  
Fluid mechanics (Navier-Stokes)