

## Green-Lagrange Strain Tensor

$$\text{right Cauchy-Green: } \underline{\underline{C}} = \underline{\underline{F}}^T \underline{\underline{F}}$$

$\Rightarrow$  quantified stretches (& shear)

$$\lambda(\hat{\underline{x}}) = |\underline{dx}| / |\underline{dx}| = \sqrt{\hat{\underline{x}} \cdot \underline{\underline{C}} \hat{\underline{x}}}$$

$$\text{No deformation: } \underline{\underline{\epsilon}} = \underline{\underline{\varphi}}(\underline{x}) = \underline{\underline{0}} \Rightarrow \underline{\underline{F}} = \underline{\underline{C}} = \underline{\underline{I}} \Rightarrow \lambda = 1$$

$$\text{Constitutive law: } \underline{\underline{\sigma}} = \underline{\underline{\sigma}}(\underline{\underline{C}})$$

no deformation = no stress

$\Rightarrow$  Strain tensor zero in absence of deformation

1D: engineering strain:  $\epsilon = \lambda - 1$

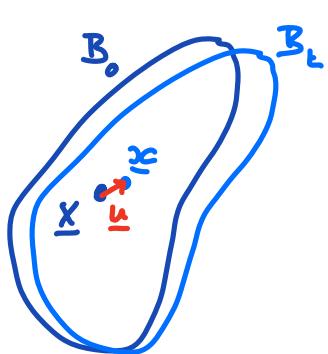
3D Green-Lagrange strain tensor

$$\underline{\underline{\epsilon}} = \frac{1}{2} (\underline{\underline{C}} - \underline{\underline{I}})$$

$$\text{no deformation: } \underline{\underline{\epsilon}} = \underline{\underline{0}} = \underline{\underline{I}} \Rightarrow \underline{\underline{\epsilon}} = \underline{\underline{0}}$$

relation to infinitesimal strain tensor

## Small displacements



natural to use displacement

$$\underline{u} = \underline{x} - \underline{X}$$

$$\underline{x} = \varphi(\underline{x})$$

$$\nabla \underline{u} = \nabla(\varphi(\underline{x}) - \underline{x}) = \nabla \varphi - \underline{\underline{I}}$$

$$\nabla \underline{u} = \underline{\underline{F}} - \underline{\underline{I}} \equiv \underline{\underline{H}}$$

Quantify magnitude of tensor:

$$|\underline{\underline{A}}| = \sqrt{\underline{\underline{A}} : \underline{\underline{A}}} = (A_{11}^2 + A_{12}^2 + \dots + A_{21}^2 + A_{22}^2)^{\frac{1}{2}}$$

Small deformation:  $|\underline{\underline{H}}| = \epsilon \ll 1$

Linearize Cauchy-Green

$$\underline{\underline{F}} = \underline{\underline{I}} + \underline{\underline{H}} \quad |\underline{\underline{H}}| = \epsilon \ll 1$$

$$\begin{aligned} \underline{\underline{C}} &= \underline{\underline{F}}^T \underline{\underline{F}} = (\underline{\underline{I}} + \underline{\underline{H}})^T (\underline{\underline{I}} + \underline{\underline{H}}) \\ &= (\underline{\underline{I}} + \underline{\underline{H}}^T)(\underline{\underline{I}} + \underline{\underline{H}}) \\ &= \underline{\underline{I}} + \underline{\underline{H}} + \underline{\underline{H}}^T + \underbrace{\underline{\underline{H}}^T \underline{\underline{H}}}_{O(\epsilon^2)} \end{aligned}$$

$$\Rightarrow \underline{\underline{C}} \approx \underline{\underline{I}} + \underline{\underline{H}} + \underline{\underline{H}}^T = \underline{\underline{I}} + \nabla \underline{u} + \nabla \underline{u}^T$$

## Linearize Euler-Lagrange

$$\underline{\underline{E}} = \frac{1}{2} (\underline{\underline{C}} - \underline{\underline{I}}) = \frac{1}{2} (\underline{\underline{I}} + \nabla \underline{\underline{u}} + \nabla \underline{\underline{u}}^T - \underline{\underline{I}}) + O(|\nabla \underline{\underline{u}}|^2)$$

$$E \approx \frac{1}{2} (\nabla \underline{\underline{u}} + \nabla \underline{\underline{u}}^T)$$

Infinitesimal strain tensor:

$$\underline{\underline{\varepsilon}} = \frac{1}{2} (\nabla \underline{\underline{u}} + \nabla \underline{\underline{u}}^T) \quad \underline{\underline{\varepsilon}} = \text{sym}(\nabla \underline{\underline{u}})$$

## Infinitesimal Stretch & Rotation

Linearize right stretch:

$$\underline{\underline{U}} = \sqrt{\underline{\underline{C}}} = (\underline{\underline{I}} + \underline{\underline{H}} + \underline{\underline{H}}^T + \underline{\underline{H}}^T \underline{\underline{H}})^{\frac{1}{2}}$$

Note: for  $\underline{\underline{A}} = \underline{\underline{A}}^T \quad m \in \mathbb{R}$

$$(\underline{\underline{I}} + \underline{\underline{A}})^m = \underline{\underline{I}} + m \underline{\underline{A}} + O(|\underline{\underline{A}}|^2)$$

shown with Taylor expansion in principal frame

$$\underline{\underline{A}} = \underline{\underline{H}} + \underline{\underline{H}}^T + O(|\underline{\underline{H}}|^2)$$

$$\Rightarrow \underline{\underline{U}} = \underline{\underline{I}} + \frac{1}{2} (\underline{\underline{H}} + \underline{\underline{H}}^T) + O(|\underline{\underline{H}}|^2)$$

$$\approx \underline{\underline{I}} + \frac{1}{2} (\nabla \underline{\underline{u}} + \nabla \underline{\underline{u}}^T) = \underline{\underline{I}} + \underline{\underline{\varepsilon}}$$

$$\text{Similarly } \underline{\underline{V}} = \sqrt{\underline{\underline{F}}\underline{\underline{F}}^T} = \underline{\underline{I}} + \frac{1}{2}(\nabla\underline{\underline{u}} + \nabla\underline{\underline{u}}^T) = \underline{\underline{I}} + \underline{\underline{\xi}}$$

Note:  $\underline{\underline{C}} \approx \underline{\underline{I}} + \underline{\underline{\xi}}$ ,  $\underline{\underline{U}} \approx \underline{\underline{I}} + \underline{\underline{\xi}}$ ,  $\underline{\underline{V}} \approx \underline{\underline{I}} + \underline{\underline{\xi}}$   
 $\Rightarrow$  linearization is the same

### Linearize rotation

$$\begin{aligned}\underline{\underline{R}} &= \underline{\underline{\Xi}} \underline{\underline{U}}^{-1} = (\underline{\underline{I}} + \underline{\underline{H}})(\underline{\underline{I}} + \underline{\underline{\xi}})^{-1} \quad \underline{\underline{\xi}} = O(|\underline{\underline{H}}|) \\ &= (\underline{\underline{I}} + \underline{\underline{H}})(\underline{\underline{I}} - \underline{\underline{\xi}}) + O(|\underline{\underline{H}}|^2) \\ &= \underline{\underline{I}} - \frac{1}{2}(\underline{\underline{H}} + \underline{\underline{H}}^T) + \underline{\underline{H}} + O(|\underline{\underline{H}}|^2) \\ &= \underline{\underline{I}} + \underline{\underline{H}} - \underline{\underline{\xi}} + O(|\underline{\underline{H}}|^2) \\ &= \underline{\underline{I}} + \frac{1}{2}(\underline{\underline{H}} - \underline{\underline{H}}^T) + O(|\underline{\underline{H}}|^2)\end{aligned}$$

$$\begin{aligned}\underline{\underline{R}} &\approx \underline{\underline{I}} + \frac{1}{2}(\nabla\underline{\underline{u}} - \nabla\underline{\underline{u}}^T) \\ &= \underline{\underline{I}} + \underline{\underline{\omega}}\end{aligned}$$

Infinitesimal Rotation Tensor:  $\underline{\underline{\omega}} = \frac{1}{2}(\nabla\underline{\underline{u}} - \nabla\underline{\underline{u}}^T)$

$\underline{\underline{\omega}} = \text{skew}(\nabla\underline{\underline{u}}) = \text{axial tensor}$

axis of rotation:  $a_j = \frac{1}{2} \epsilon_{mijn} \omega_{mn}$

Linearization of  $\underline{F}$

$$\underline{U} \approx \underline{I} + \underline{\varepsilon} \quad \underline{R} = \underline{I} + \underline{\omega}$$

$$\underline{F} = \underline{R} \underline{U} = (\underline{I} + \underline{\varepsilon}) (\underline{I} + \underline{\omega}) = \underline{I} + \underline{\varepsilon} + \underline{\omega} + O(|H|^2)$$

$$\Rightarrow \underline{F} = \underline{R} \underline{U} \approx \underline{I} + \underline{\varepsilon} + \underline{\omega}$$

$\underbrace{\varepsilon}_{\text{multiplicative}}$      $\underbrace{\omega}_{\text{additive}}$

$$\underline{F} = \underline{I} + \nabla \underline{u} = \underline{I} + \text{sym}(\nabla \underline{u}) + \text{skew}(\nabla \underline{u})$$

$\Rightarrow$  nature of the rotation-stretch decomposition  
changes from multiplicative to additive!

## Interpretation of components of $\underline{\underline{\epsilon}}$

Start from  $\underline{\underline{\epsilon}} = \underline{\underline{\epsilon}}_0 + \underline{\underline{\epsilon}}_e$

$$\underline{\underline{\epsilon}} = \underline{\underline{\epsilon}}_0 = \frac{1}{2} (\underline{\underline{\epsilon}} - \underline{\underline{\epsilon}}_e) + O(\epsilon^2)$$

$$|\nabla u| = \epsilon$$

$\uparrow$   
scalar

$$\Rightarrow \underline{\underline{\epsilon}} \approx \underline{\underline{\epsilon}}_e + 2 \underline{\underline{\epsilon}}_0$$

I) Diagonal components

$$C_{ii} = 1 + 2 \epsilon_{ii}$$

$$\sqrt{C_{ii}} = \lambda(\epsilon_i)$$

$$\lambda(\epsilon_i) = \sqrt{1 + 2 \epsilon_{ii}}$$

Expand in Taylor-Series:  $\sqrt{1+x} = 1 + \frac{x}{2} - \frac{x^2}{8} + \dots$

where  $x = 2 \epsilon_{ii}$  s.t.  $\lambda(\epsilon_i) = 1 + \epsilon_{ii}$

$\Rightarrow \boxed{\epsilon_{ii} \approx \lambda(\epsilon_i) - 1}$  engineering strain  
in coordinate directions

$$\lambda = \frac{|y - z|}{|x - z|} = \frac{l}{L} \text{ stretch}$$

$$\lambda - 1 = \frac{|y - z| - |x - z|}{|x - z|} = \frac{l - L}{L} = \frac{\Delta l}{L}$$

$\epsilon_{ii}$  = relative change in length

## II) Off-Diagonal Components

$$C_{ij} = \lambda(\epsilon_i) \lambda(\epsilon_j) \sin(\gamma_{ij}) \quad \gamma_{ij} = \gamma(\epsilon_i, \epsilon_j)$$

$$C_{ij} \approx 2 \epsilon_{ij} \quad i \neq j$$

$$\epsilon_{ij} \approx \frac{1}{2} \lambda(\epsilon_i) \lambda(\epsilon_j) \sin(\gamma_{ij})$$

$$\lambda(\epsilon_i) = 1 + \epsilon_{ii} = 1 + O(\epsilon)$$

$$\lambda(\epsilon_j) = 1 + \epsilon_{jj} = 1 + O(\epsilon)$$

$$\sin(\gamma_{ij}) = O(\epsilon)$$

$$\epsilon_{ij} = \frac{1}{2} (1 + \epsilon_{ii}) (1 + \epsilon_{jj}) \sin(\gamma_{ij})$$

$$= \frac{1}{2} \left[ \sin(\gamma_{ij}) + \underbrace{\epsilon_{ii} \sin(\gamma_{ij})}_{O(\epsilon^2)} + \underbrace{\epsilon_{jj} \sin(\gamma_{ij})}_{O(\epsilon^2)} + \underbrace{\epsilon_{ii} \epsilon_{jj} \sin(\gamma_{ij})}_{O(\epsilon^3)} \right]$$

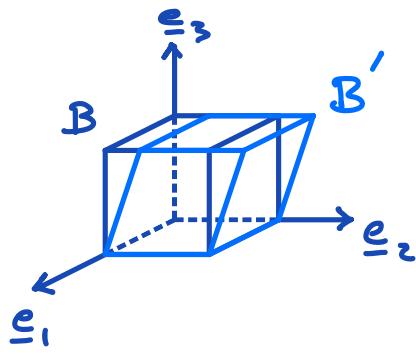
$$\Rightarrow \epsilon_{ij} = \frac{1}{2} \sin(\gamma(\epsilon_i, \epsilon_j))$$

In infinitesimal strain  $\Rightarrow \gamma \ll 1 \Rightarrow \sin \gamma \approx \gamma$

$$\boxed{\epsilon_{ij} \approx \frac{1}{2} \gamma(\epsilon_i, \epsilon_j)}$$

$\Rightarrow$  half shear angle between coord. directions

## Example: Simple shear



$$\underline{\underline{\epsilon}} = \underline{\varphi}(\underline{X}) = \begin{bmatrix} \varphi_1 \\ \varphi_2 \\ \varphi_3 \end{bmatrix} = \begin{bmatrix} X_1 \\ X_2 + \alpha X_3 \\ X_3 \end{bmatrix} \quad \alpha > 0$$

large def:  $\underline{\underline{\epsilon}} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & \alpha \\ 0 & \alpha & 1+\alpha^2 \end{bmatrix}$

Infinitesimal:

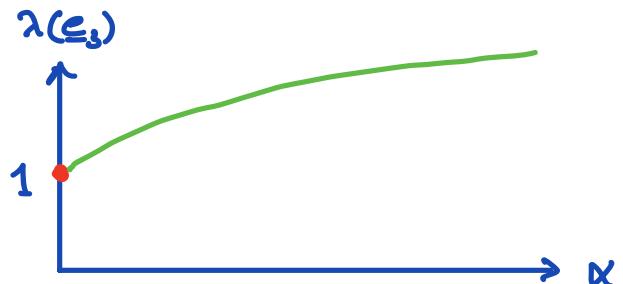
$$\underline{u} = \underline{\underline{\epsilon}} \underline{e} - \underline{X} = \begin{bmatrix} 0 \\ \alpha X_3 \\ 0 \end{bmatrix} \quad \nabla \underline{u} = \underline{\underline{H}} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & \alpha \\ 0 & \alpha & 0 \end{bmatrix}$$

$$\underline{\underline{\epsilon}} = \frac{1}{2} (\nabla \underline{u} + \nabla \underline{u}^T) = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & \alpha/2 \\ 0 & \alpha/2 & 0 \end{bmatrix}$$

Infinitesimal stretches:  $\epsilon_{ii} = \lambda(e_i) - 1 = 0$

$\lambda(e_i) = 1$  no stretch in any coord. direction

Finite stretches:  $C_{ii} = \lambda^2(e_i)$   $C_{33} = \lambda^2(e_3) = 1 + \alpha^2$



$$\lambda(e_3) = \sqrt{1 + \alpha^2}$$

- finite strain  
- infinitesimal strain

Infinitesimal shear :  $\varepsilon_{ij} \approx \frac{1}{2} \gamma(\varepsilon_i, \varepsilon_j)$

$$\gamma(\varepsilon_2, \varepsilon_3) = 2 \varepsilon_{23} = \alpha$$

Finite shear :  $\gamma(\varepsilon_2, \varepsilon_3) = \alpha \sin\left(\frac{\alpha}{\sqrt{1+\alpha^2}}\right)$

$$\gamma(\varepsilon_2, \varepsilon_3)$$



- finite strain

- infinitesimal