

# Discretizing the Stokes Equation

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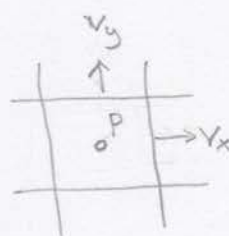
Dimensionless variable viscosity Stokes equation:

$$\begin{array}{l} 1) \quad \nabla \cdot [\mu(\nabla \underline{v} + \nabla^T \underline{v})] - \nabla p = \underline{f}_1 \\ 2) \quad \nabla \cdot \underline{v} = \underline{f}_2 \end{array}$$

Overall we are looking for a discrete system of the form

$$\begin{bmatrix} \underline{A} & \underline{C}^T \\ \underline{C} & 0 \end{bmatrix} \begin{bmatrix} \underline{v} \\ p \end{bmatrix} = \begin{bmatrix} \underline{f}_1 \\ \underline{f}_2 \end{bmatrix} \quad \text{where } \underline{A} = \underline{A}^T \Rightarrow \text{system is symmetric}$$

Here  $\underline{v} = \begin{bmatrix} v_x \\ v_y \end{bmatrix}$  on cell faces and  $p$  is a vector of cell center pressures.



First some low-hanging fruit:

From the discrete system

$$1) \quad \underline{A} * \underline{v} + \underline{C}^T * p = \underline{f}_1$$

$$2) \quad \underline{C} * \underline{v} = \underline{f}_2$$

From 2:

$$\nabla \cdot \underline{v} \approx \underline{C} * \underline{v} \Rightarrow \underline{C} = \underline{D} \quad \text{standard discrete divergence on our normal staggered grid}$$

From 1:

$$-\nabla p \approx \underline{C}^T p \Rightarrow \underline{C}^T = -\underline{G} \quad \text{standard discrete gradient} \\ = -(-\underline{D}^T) = \underline{C}^T$$

$\Rightarrow$  build on the existing staggered grid because it gets the relations between  $p$  and  $\underline{v}$  right.

## Divergence of deviatoric stress tensor

(2)

$$\text{Cauchy stress: } \underline{\underline{\sigma}} = \underbrace{-p \underline{\underline{I}}}_{\text{vol.}} + \underbrace{2\mu \underline{\underline{\dot{\epsilon}}}}_{\text{deviatoric stress}} = -p \underline{\underline{I}} + \underline{\underline{\tau}} \quad (\underline{\underline{\tau}} = 2\mu \underline{\underline{\dot{\epsilon}}}) = -p \underline{\underline{I}} + \underline{\underline{\tau}}$$

$$\text{Deviatoric stress: } \underline{\underline{\tau}} = 2\mu \underline{\underline{\dot{\epsilon}}} = \mu (\nabla \underline{\underline{v}} + \nabla^T \underline{\underline{v}})$$

$$\nabla \cdot \underline{\underline{\sigma}} = \nabla \cdot (\underline{\underline{\tau}} - p \underline{\underline{I}}) = \nabla \cdot \underline{\underline{\tau}} - \underbrace{\nabla p}_{\text{done}}$$

⇒ need to discretize divergence of deviatoric stress tensor.

Definition of the divergence of a 2nd order tensor:

$$\nabla \cdot \underline{\underline{\sigma}} = \sigma_{ij,j} \hat{e}_i = \begin{pmatrix} \sigma_{11,1} & \sigma_{12,2} \\ \sigma_{21,1} & \sigma_{22,2} \end{pmatrix} = \begin{pmatrix} \nabla \cdot (\sigma_{11} \sigma_{12}) \\ \nabla \cdot (\sigma_{21} \sigma_{22}) \end{pmatrix}$$

⇒ Divergence is applied row-wise to the tensor

1<sup>st</sup>-row is the divergence of the x-stresses  
2<sup>nd</sup>-row is the divergence of the y-stresses

Definition of the gradient of a vector:

$$\nabla \underline{\underline{v}} = v_{ij} \hat{e}_i \otimes \hat{e}_j = \begin{pmatrix} v_{1,1} & v_{2,1} \\ v_{2,1} & v_{2,2} \end{pmatrix} = \begin{pmatrix} \nabla v_1 \\ \nabla v_2 \end{pmatrix} \quad \text{hence } \nabla^T \underline{\underline{v}} = (\nabla \underline{\underline{v}})^T = (\nabla v_1, \nabla v_2)$$

⇒ Gradient is also applied row-wise

Hence the symmetric velocity gradient/strain-rate tensor

$$\underline{\underline{\dot{\epsilon}}} = \frac{1}{2} (\nabla \underline{\underline{v}} + \nabla^T \underline{\underline{v}}) = \begin{pmatrix} v_{1,1} & \frac{1}{2}(v_{1,2} + v_{2,1}) \\ \frac{1}{2}(v_{2,1} + v_{1,2}) & v_{2,2} \end{pmatrix} \Rightarrow \boxed{\underline{\underline{\dot{\epsilon}}} = \underline{\underline{\dot{\epsilon}}}^T} \quad \text{symmetric}$$

# Discretizing the strain-rate tensor

(3)

To discretize the strain-rate tensor we need the following derivatives

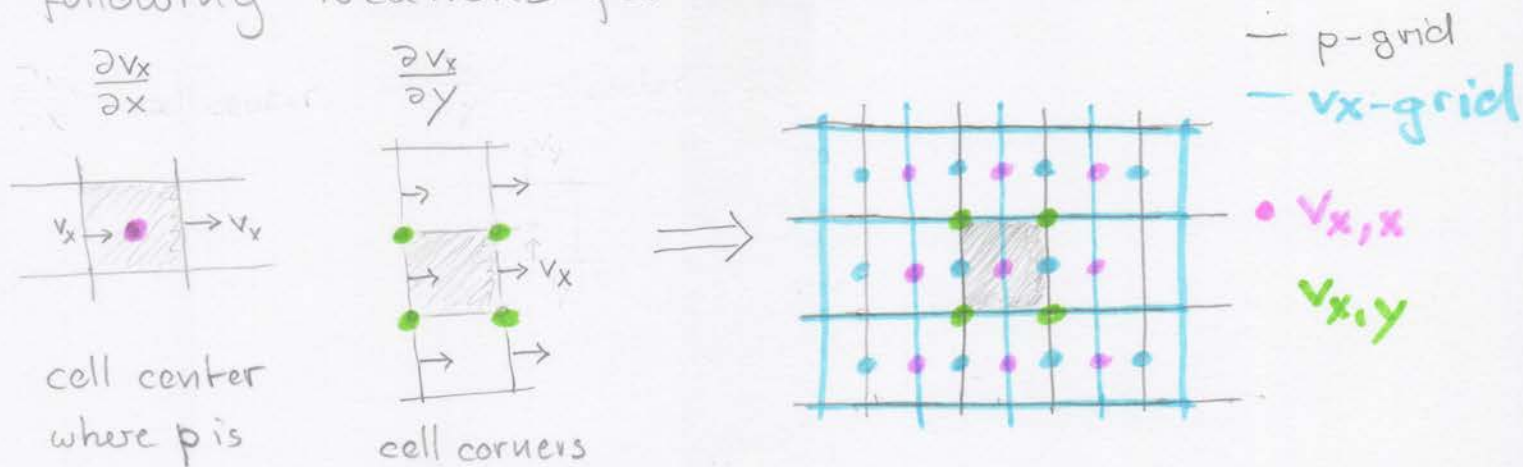
$$v_{1,1} = \frac{\partial v_1}{\partial x_1} = \frac{\partial v_x}{\partial x}$$

$$v_{1,2} = \frac{\partial v_1}{\partial x_2} = \frac{\partial v_x}{\partial y}$$

$$v_{2,2} = \frac{\partial v_2}{\partial x_2} = \frac{\partial v_y}{\partial y}$$

$$v_{2,1} = \frac{\partial v_2}{\partial x_1} = \frac{\partial v_y}{\partial x}$$

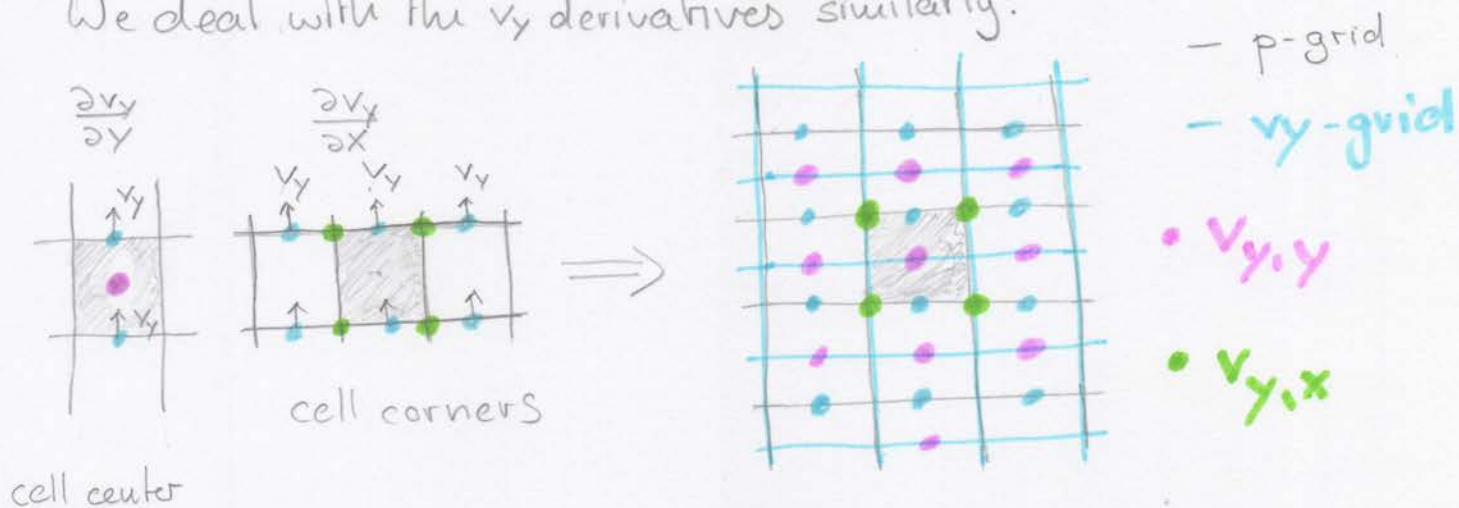
Given the standard staggered grid we have the following locations for these derivatives.



⇒ Simplest way to compute the  $v_x$ -derivatives is to introduce a new grid centered on  $v_x$ .

This  $v_x$ -grid is shifted by  $\frac{\Delta x}{2}$  relative to the p-grid and its size is  $N_x+1$  by  $N_y$  if  $N_x$  and  $N_y$  are the size of the pressure grid.

We deal with the  $v_y$  derivatives similarly:



## 2D Stokes grid

(4)

• In 2D we use 3 staggered grids

1) Pressure grid: Primary grid that defines location of pressure and velocities  
 $N_x$  by  $N_y$

⇒ used for heat transport calculations

2) X velocity grid: Shifted by  $\frac{\Delta x}{2}$  in x-dir relative to pressure grid. Used to compute  $v_x$  derivatives in strain-rate tensor  
 $N_x + 1$  by  $N_y$

3) Y velocity grid: Shifted by  $\frac{\Delta y}{2}$  in y-dir relative to pressure grid. Used to compute  $v_y$  derivatives in strain-rate tensor.  
 $N_x$  by  $N_y + 1$

Note: In 3D we would have one additional grid for  $v_z$  velocity.

⇒ lots of book keeping.

New Matlab function:

`Grid = build_stokes_grid(Gridp)`

That builds all other grids given the primary pressure grid and returns a structure that holds all grids.

`Grid.p` = pressure grid

`Grid.x` = x-velocity grid

`Grid.y` = y-velocity grid

} All three grids are built using standard `buildgrid.m` we have already developed

⇒ this will be on IfW 9