

Conductive cooling of a finite domain

```
clear, close all, clc
set_demo_defaults()
```

Dimensional Problem

To determine the cooling timescale of a finite domain of length, L , with initial temperature, T_0 , and boundary temperature, T_b , we need to solve the following Initial and Boundary Value Problem (IBVP)

$$\text{PDE: } \rho c_p \frac{\partial T}{\partial t} - \nabla \cdot [\kappa \nabla T] = 0 \quad \text{on } x \in [0, L] \text{ and } t \geq 0.$$

$$\text{BC's: } T(x=0, t) = T_b \quad \text{and} \quad \nabla T \cdot \hat{\mathbf{n}}|_{x=L} = 0.$$

$$\text{IC: } T(x, t=0) = T_0 \quad \text{on } x \in [0, L]$$

Here we assume that c_p and κ are constant so that the problem is linear. Note that we assume that $T_0 > T_b$ so that $\Delta T = T_0 - T_b > 0$.

This problem has six parameters: L , c_p , ρ , κ , T_0 and T_b

To determine how many independent parameters this problem has we non-dimensionalize the variables, T , x and t by choosing the following characteristic scales.

$$T' = \frac{T - T_b}{\Delta T}, \quad x' = \frac{x}{L} \quad \text{and} \quad t' = \frac{t}{t_c}.$$

Here we have chosen external scales imposed by the domain size and boundary for T and x . There is no external scale for t and here we just use a generic characteristic time, t_c , that will be determined by the equation itself.

Dimensionless Problem

Substituting the generic scales we obtain the following problem

$$\text{PDE: } \frac{\partial T'}{\partial t'} = \Pi \nabla'^2 T' \quad \text{on } x' \in [0, 1] \text{ and } t' \geq 0.$$

$$\text{BC's: } T'(0, t') = 0 \quad \text{and} \quad \nabla' T' \cdot \hat{\mathbf{n}}|_{x'=1} = 0.$$

$$\text{IC: } T'(x', 0) = 1 \quad \text{on } x' \in [0, 1]$$

where the dimensionless group $\Pi = \frac{\kappa t_c}{\rho c_p L^2} = \frac{t_c}{t_D}$ is the ratio of t_c and the internal diffusive

timescale $t_D = \frac{\rho c_p L^2}{\kappa} = \frac{L^2}{D}$, where $D = \frac{\kappa}{\rho c_p}$ is the thermal diffusivity. To simplify the problem we choose

the internal diffusive timescale as characteristic scale $t_c = t_D$. So that the final dimensionless problem is parameter free and here written in one dimension as

$$\text{PDE: } \frac{\partial T'}{\partial t'} = \frac{\partial^2 T'}{\partial x'^2} \quad \text{on } x' \in [0, 1] \text{ and } t' \geq 0.$$

$$\text{BC's: } T'(0, t') = 0 \quad \text{and} \quad \left. \frac{\partial T'}{\partial x'} \right|_{x'=1} = 0.$$

$$\text{IC: } T'(x', 0) = 1 \quad \text{on } x \in [0, 1]$$

Analytic solution

The analytic solution is obtained by separation of variables, $T'(x', t') = g(t')h(x')$ and the dimensionless solution is given by

$$T'(x', t') = \frac{4}{\pi} \sum_{n=1}^{\infty} \frac{1}{2n-1} \sin\left(\left(n - \frac{1}{2}\right)\pi x'\right) \exp\left(-\left(n - \frac{1}{2}\right)^2 \pi^2 t'\right)$$

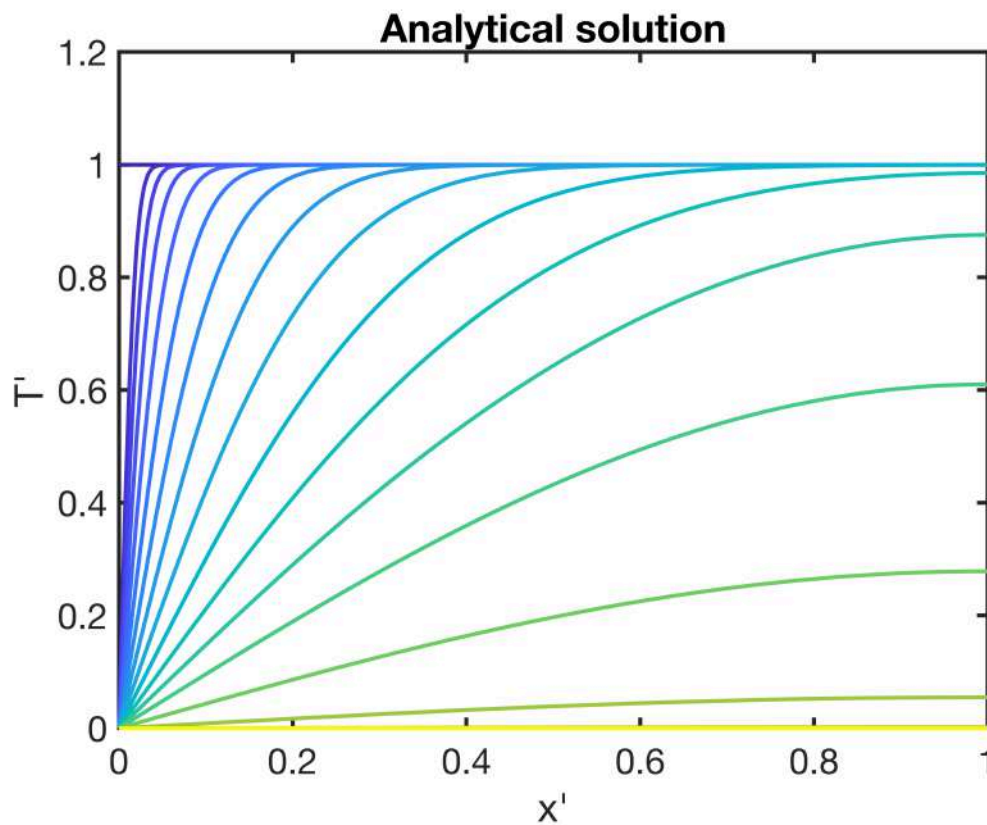
The corresponding dimensional solution is obtained by substituting the definitions of the dimensionless variables and given by

$$T(x, t) = T_b + \frac{4\Delta T}{\pi} \sum_{n=1}^{\infty} \frac{1}{2n-1} \sin\left(\frac{\left(n - \frac{1}{2}\right)\pi x}{L}\right) \exp\left(-\frac{\left(n - \frac{1}{2}\right)^2 \pi^2 D}{L^2} t\right)$$

For simplicity we work with the dimensionless solution below.

```
Nt = 20; Nx = 1e3; Nn = 1e2;

tvec = logspace(-4, 2, Nt);
figure
map = colormap(parula(Nt+1));
plot([0 1], [1 1], 'color', map(1, :)), hold on
for i = 1:Nt
    [xd, Td] = analytic_soln_finite(tvec(i), Nx, Nn);
    plot(xd, Td, 'color', map(i+1, :))
end
title 'Analytical solution'
xlabel 'x'', ylabel 'T''
```



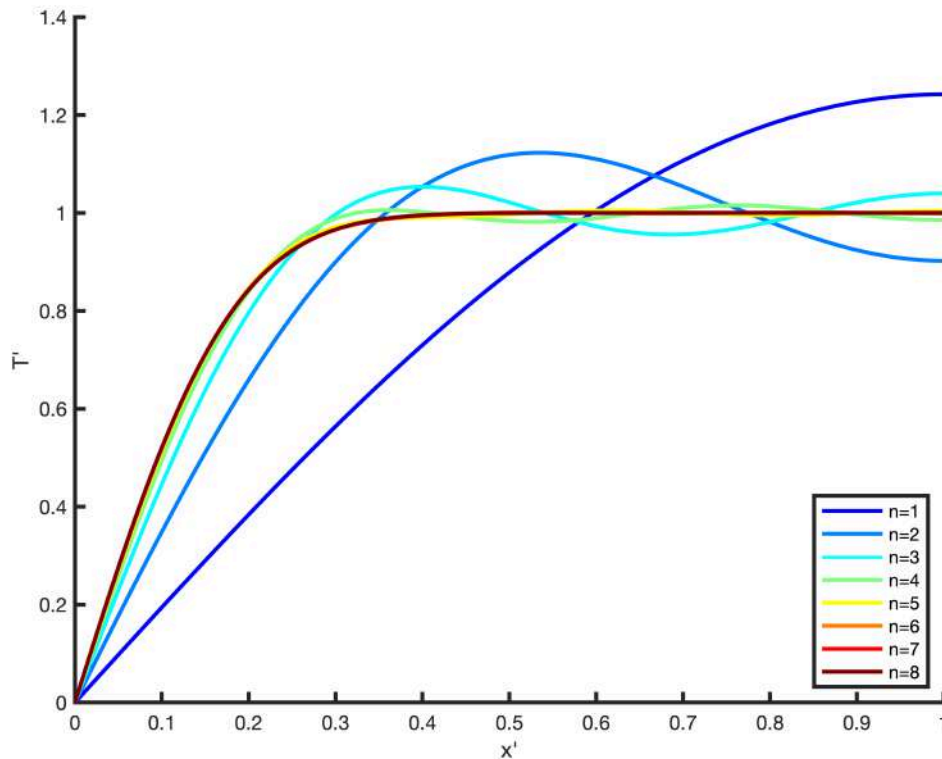
drawnow

Convergence of the series

The analytic solution is a sum of sin functions of increasing wave number that are superposed to generate the solution

```
Nx = 1e3; Nn = 8;

figure
map = colormap(jet(Nn));
for i = 1:Nn
    [xd,Td] = analytic_soln_finite(.01,Nx,i); hold on
    plot(xd,Td,'color',map(i,:))
end
legend('n=1','n=2','n=3','n=4','n=5','n=6','n=7','n=8','location','southeast')
xlabel 'x'', ylabel 'T''
```



drawnow

Temperature decay

The contribution of each eigenfunction/mode decays proportional to its eigenvalue. The higher order modes decay quickly and at late time the solution is well approximated by the first mode.

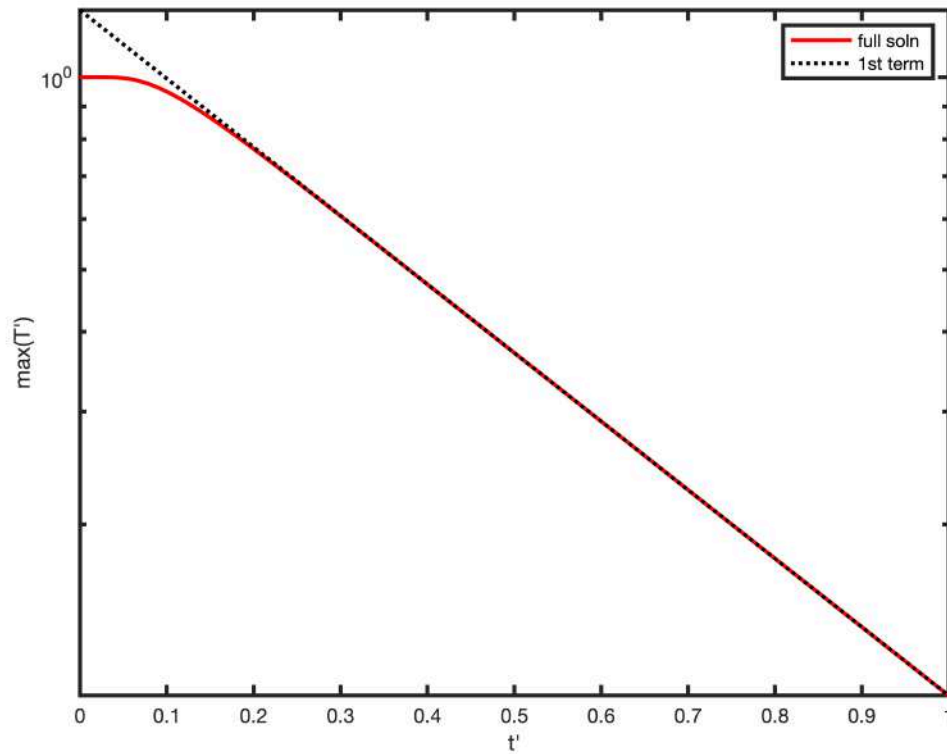
$$T'(x', t') \approx \frac{4}{\pi} \sin\left(\frac{\pi x'}{2}\right) \exp\left(-\frac{\pi^2 t'}{4}\right)$$

The figure show the decay of the maximum temperature from the full solution and the approximation

```
Nt = 200; Nx = 2;
tvec = linspace(0,1,Nt);
Td_max = zeros(Nt,1); Td_max_approx = Td_max;
figure
for i = 1:Nt
    [xd,Td] = analytic_soln_finite(tvec(i),Nx,1e3);
    Td_max(i) = Td(2);
    [xd,Td] = analytic_soln_finite(tvec(i),Nx,1);
    Td_max_approx(i) = Td(2);
end
semilogy(tvec,Td_max,'r-'), hold on
semilogy(tvec,Td_max_approx,'k:')

ylabel 'max(T')', xlabel 't''
```

```
legend('full soln','1st term')
```



```
drawnow
```

Heat flow (to be continued)

The heat flow at the boundary is given by Fourier's law in its dimensionless form

$$q = -\kappa \frac{\partial T}{\partial x}$$

or in dimensionless form

$$q' = -\frac{\partial T'}{\partial x'}$$

if we introduce the characteristic scale, $q_c = K \frac{\Delta T}{L}$ for the heat flow. Differentiating the analytic solution we have

Numerical solution

Below is the code to solve the dimensionless problem numerically

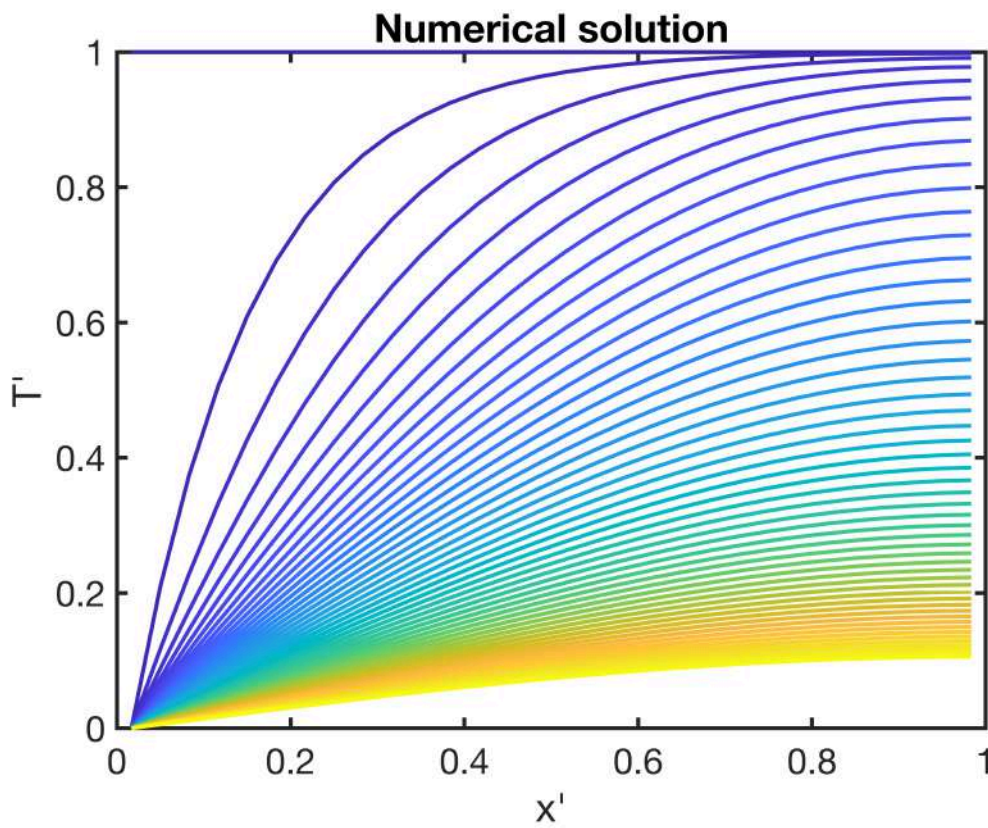
```
tmax = 1;
Nt = 50;
theta = 0;

% Grid and operators
Grid.xmin = 0; Grid.xmax = 1; Grid.Nx = 30;
Grid = build_grid(Grid);
[D,G,~,I,M] = build_ops(Grid);
L = -D*G; fs = zeros(Grid.Nx,1);
IM = @(theta,dt) I + (1-theta)*dt*L;
EX = @(theta,dt) I - theta*dt*L;

% Boundary conditions
BC.dof_dir = Grid.dof_xmin;
BC.dof_f_dir = Grid.dof_f_xmin;
BC.g = 0;
BC.dof_neu = [];
BC.dof_f_neu = [];
BC.qb = [];
[B,N,fn] = build_bnd(BC,Grid,I);

% Initial condition
u = ones(Grid.Nx,1);
figure
map = colormap(parula(Nt+1));
plot(Grid.xc,u,'color',map(1,:)), hold on

% Timestepping
dt_max = Grid.dx^2/2;
dt = tmax/Nt;
time = zeros(Nt+1,1);
for n = 1:Nt
    time(n+1) = time(n) + dt;
    u = solve_lbvp(IM(theta,dt),dt*fs+EX(theta,dt)*u,B,BC.g,N);
    plot(Grid.xc,u,'-', 'color',map(n+1,:)), drawnow
end
title 'Numerical solution'
xlabel 'x'', ylabel 'T''
```



```
function [xd,Td] = analytic_soln_finite(td,Nx,Nn)
% Input:
% td = dimensionless time
% Nx = number of gridpoints
% Nn = number of terms in the Fourier series

% Output:
% xd = dimensionless distance
% Td = dimensionless temperature
[xd] = linspace(0,1,Nx)';
Td = zeros(Nx,1);
for n = 1:Nn
    Td = Td + 1/(2*n-1)*sin((n-.5)*pi*xd)*exp(-(n-.5).^2*pi^2*td);
end
Td = 4/pi*Td;
end
```