

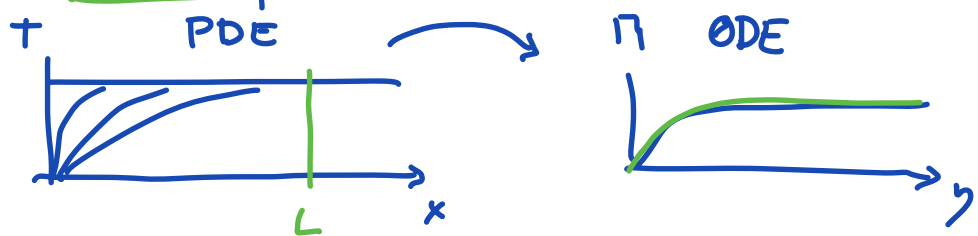
Lecture 20: Advection

Logistics: - HW8 is due (8/11)

⇒ office hours

Last time: - Transient heat conduction

semi-infinite domain



- similarity variable:

$$\eta = \frac{x}{\sqrt{4Dt}}$$

⇒ $x \sim \sqrt{t}$ heat propagation

Today: - Advection Equation
- Method of characteristics
- Discretization

Advection

Lecture 15: Energy balance

$$\underbrace{\bar{\rho} c_p}_{(1)} \frac{\partial T}{\partial t} + \nabla \cdot \left[\underbrace{q + \rho_f c_p f T}_{(2)} - \underbrace{\bar{\kappa} \nabla T}_{\text{GW flow}} \right] = 0$$

Lectures 16-19:

Conduction $\underbrace{\rho c_p}_{\text{transfer}} \frac{\partial T}{\partial t} - \nabla \cdot \kappa \nabla T = \rho H$

If $\bar{\rho} c_p = \text{const.}$

$$\frac{\partial T}{\partial t} + \nabla \cdot [\underline{v}_e T - \bar{\alpha} \nabla T] = 0$$

Advection
Diff. Equ.

$$\underline{v}_e = \frac{\phi \rho_f c_p f}{\bar{\rho} c_p} \underline{v}_f \quad q = \phi \underline{v}_f$$

$$\bar{\alpha} = \frac{\bar{\kappa}}{\bar{\rho} c_p} \quad \text{thermal diffusivity}$$

Scaling:

$$T' = \frac{T}{T_c}$$
$$\underline{v}' = \frac{\underline{v}_e}{|\underline{v}_e|}$$

$$x' = \frac{x}{x_c}$$
$$L$$

$$t' = \frac{t}{t_c} = t_A$$

typically

T_c and

x_c are given (external)

$$\frac{\cancel{T_c}}{t_c} \frac{\partial T'}{\partial t'} + \frac{1}{x_c} \nabla' \cdot \left[\underline{v_e} \cancel{T_c} T' - \frac{\bar{\alpha} \cancel{T_c}}{x_c} \nabla' T' \right] = 0$$

$$\frac{\partial T'}{\partial t'} + \nabla' \cdot \left[\underbrace{\frac{|v_e| t_c}{x_c}}_{\underline{\pi_1}} \nabla' T' - \underbrace{\frac{\bar{\alpha} t_c}{x_c^2}}_{\underline{\pi_2}} \nabla' T' \right] = 0$$

$$\pi_1 = \frac{|v_e| t_c}{x_c} = 1 \Rightarrow t_c = \frac{x_c}{|v_e|} = t_A \text{ advective time scale}$$

$$\pi_2 = \frac{\bar{\alpha} t_c}{x_c^2} = 1 \Rightarrow t_c = \frac{x_c^2}{\bar{\alpha}} = t_D \text{ diffusive time scale}$$

$\left(\frac{L^2}{D} \right)$
 $x \sim \sqrt{t}$

choose: $t_c = t_A$

$$\frac{\partial T'}{\partial t'} + \nabla' \cdot \left[\underline{v'} T' - \left(\frac{t_A}{t_D} \right) \nabla' T' \right]$$

$$\frac{t_A}{t_D} = \frac{x_c \bar{\alpha}}{v_e x_c^2} = \frac{\bar{\alpha}}{v_e x_c} = \frac{1}{Pe}$$

$$Pe = \left(\frac{v L}{D} \right) = \frac{v_e x_c}{\bar{\alpha}} = \frac{t_D}{t_A}$$

↑
 muss
 notiertu

Peclet number

$$\frac{\partial T'}{\partial t'} + \nabla' \cdot [\underline{v}' T' - \left(\frac{1}{Pe}\right) \nabla' T'] = 0$$

Look at $Pe \gg 1$

$$\Rightarrow \frac{\partial T'}{\partial t'} + \nabla' \cdot [\underline{v}' T'] = 0$$

Re dimensionality

$$\boxed{\frac{\partial T}{\partial t} + \nabla \cdot [\underline{v}_e T]} = 0 \quad \text{Advection equation}$$

Side note : $t_e = t_D$

$$= \frac{\partial T'}{\partial t'} + \nabla' \cdot [Pe \underline{v}' T' - \nabla' T'] = 0$$

$$Pe = \frac{t_D}{t_r}$$

Solution to the Advection equation

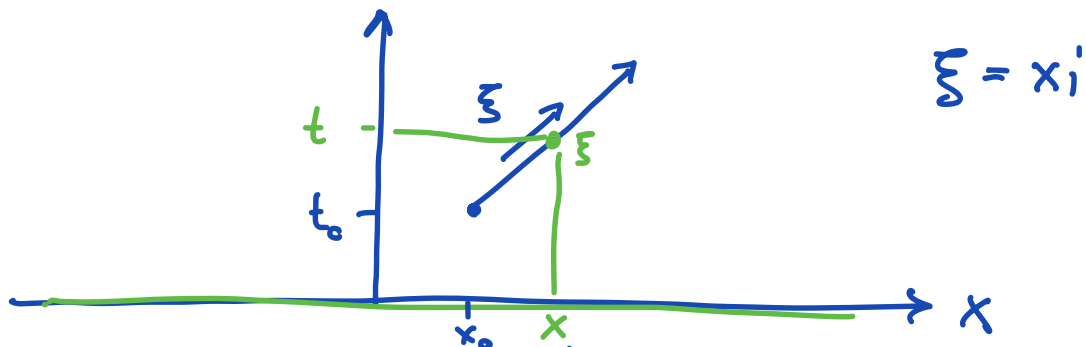
Consider heat transport in 1D column
with const. $v_e = v = \text{const}$

$$\begin{aligned} \nabla \cdot [\underline{v}_e T] &= \underline{v}_e \cdot \nabla T + T (\nabla \cdot \underline{v}_e) \\ &= \underline{v}_e \cdot \nabla T \stackrel{1D}{=} v_e \frac{\partial T}{\partial x} \end{aligned}$$

$$\text{PDE: } \frac{\partial T}{\partial t} + v \frac{\partial T}{\partial x} = 0 \quad x \in \mathbb{R}$$

$$\text{IC: } T(x, 0) = T_0(x)$$

Solve with Method of Characteristics



Idea: Find a characteristic curve, ξ ,
along which PDE reduces to ODE
 $T(x, t) = T(x(\xi), t(\xi)) = \Theta(\xi)$

Total change in T :

$$\frac{d\Theta}{d\xi} = \frac{d}{d\xi} (T(x(\xi), t(\xi))) = \frac{\partial T}{\partial t} \frac{dt}{d\xi} + \frac{\partial T}{\partial x} \frac{dx}{d\xi}$$

compare with PDE: $\frac{\partial T}{\partial t} + \frac{\partial T}{\partial x} v = 0$

$$1) \frac{d\Theta}{d\xi} = 0$$

$$2) \frac{dt}{d\xi} = 1$$

$$3) \frac{dx}{d\xi} = v$$

Combine 2 & 3: $\frac{dx}{dt} = v$

Solve for characteristic: $x - x_0 = v(t - t_0)$
char. equ

At initial condition:

$$c(x = x_0, t = t_0) = c_0(x_0)$$

substitute characteristic $x_0 = x - v(t - t_0)$

General analytic solu for Advection equ

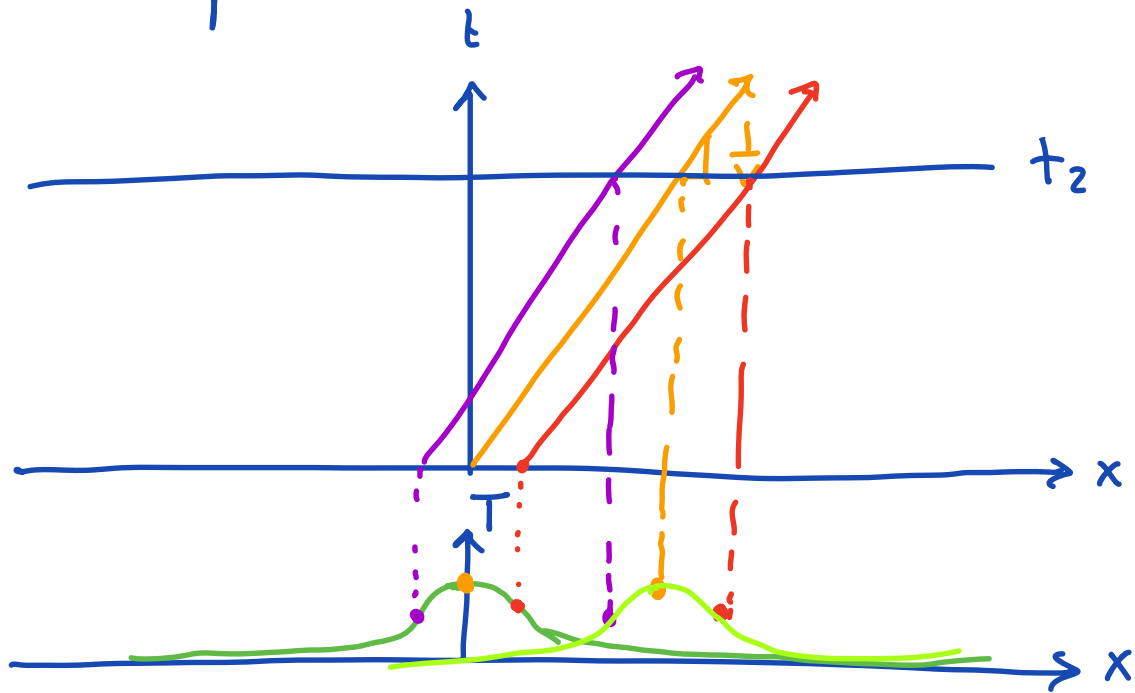
$$c(x, t) = c_0(x - v(t - t_0))$$

typically $t_0 = 0 \Rightarrow c(x, t) = c_0(\underbrace{x - vt}_{\text{travelling wave coord.}})$

Definition:

Wave is a signal/disturbance/variation moving through a medium with a recognizable speed of propagation.

⇒ shifts IC to right ($v > 0$) without change in shape



$$x - x_0 = v t$$

$$t = \frac{1}{v} (x - x_0)$$

Steady advection diffusion

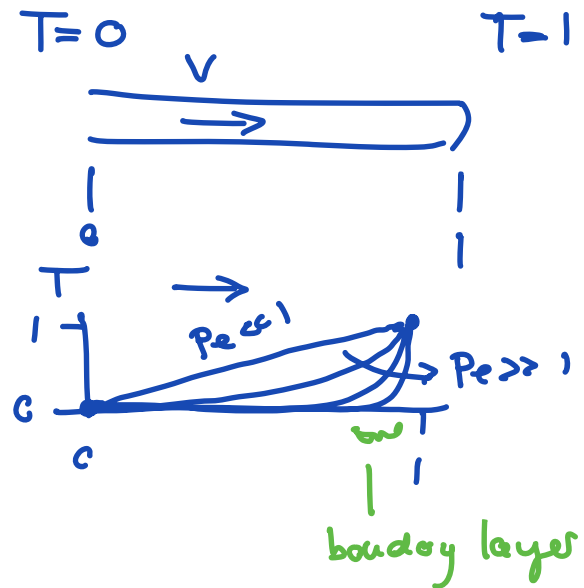
Steady flow

$$\text{PDE: } \frac{d}{dx} \left(Pe T - \frac{dT}{dx} \right) = 0$$

$$\text{BC: } T(0) = 0 \quad T(1) = 1$$

Analytic solution:

$$T(x) = \frac{e^{Pe x} - 1}{e^{Pe} - 1}$$



Discretize:

$$\text{PDE: } \nabla \cdot (Pe \underline{q} T - \nabla T) = 0$$

$$\text{Discrete: } \underline{D} * (Pe \underline{A}(\underline{q}) \underline{u} - \underline{G} \underline{u}) = 0$$

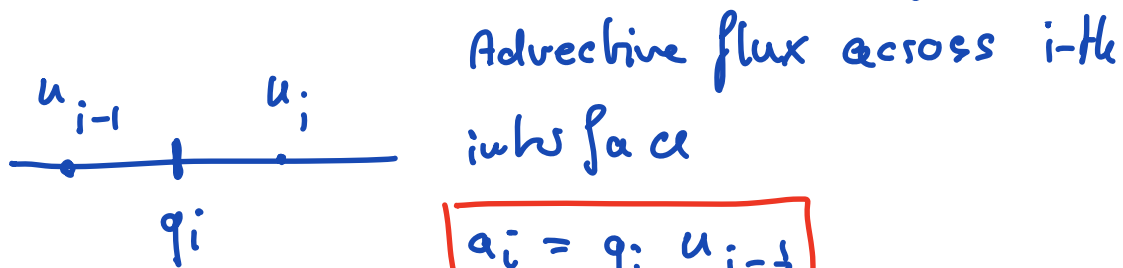
$$\underline{L} \underline{u} = 0$$

$$\underline{L} = \underline{D} * [Pe \underline{A}(\underline{q}) - \underline{G}]$$

\underline{A} is $N \times +1$ by $N \times$ matrix that is function of known fluxes \underline{q}

$$\underline{A}(q) \underline{u} \approx q^T$$

Problem is that u is given on cell centers but flux needs to be on cell faces.



$$a_i = q_i u_{i-\frac{1}{2}}$$

Need to approximate $u_{i-\frac{1}{2}}$

Central flux

$$u_{i-\frac{1}{2}} = \frac{1}{2} (u_i + u_{i-1})$$

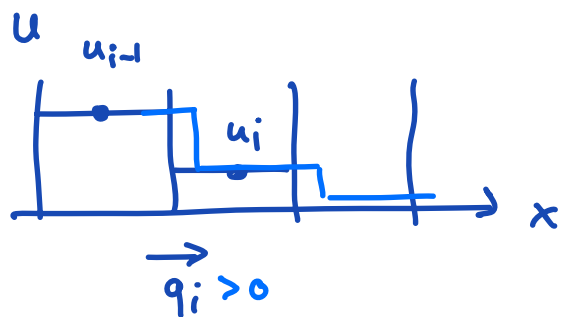
\Rightarrow leads to oscillations

Upwind flux

from HOC we know that

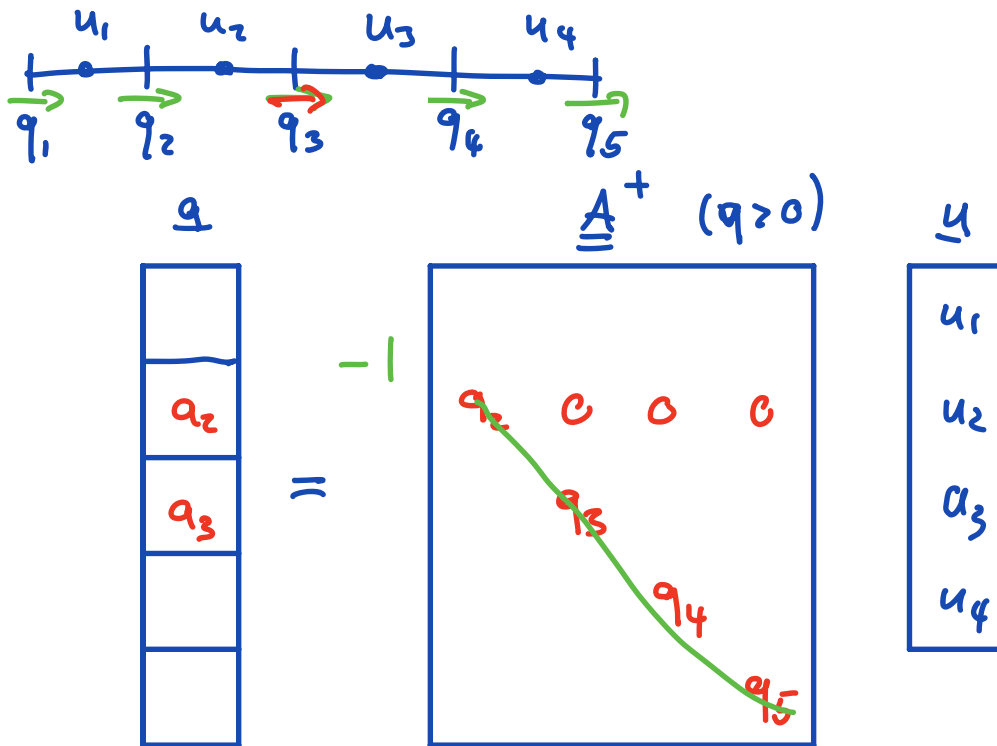
solution is simply shifted

to the right ($q \geq 0$). $\Rightarrow u_{i-\frac{1}{2}} =$



$$u_{i-\frac{1}{2}} = \begin{cases} u_{i-1} & \text{if } q_i > 0 \\ u_i & \text{if } q_i < 0 \end{cases}$$

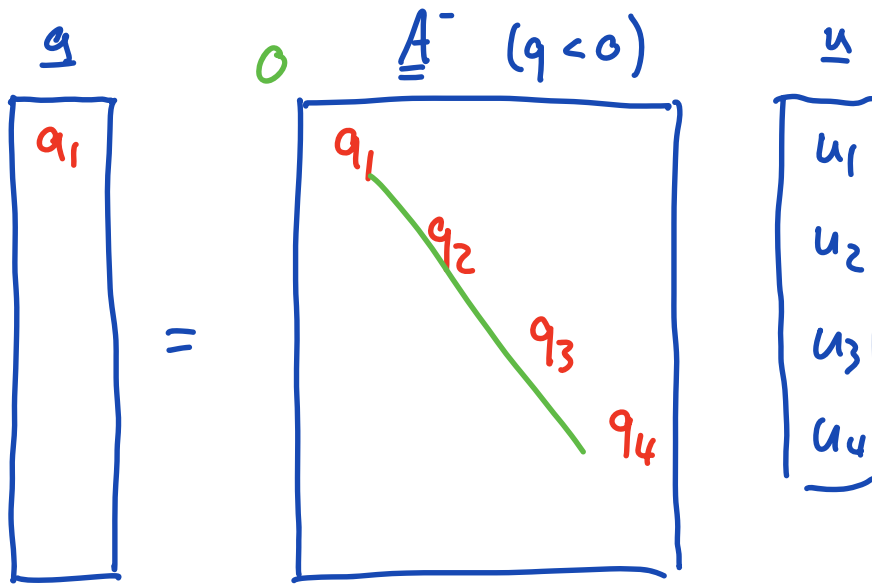
What is \underline{A} :



$$a_2 = q_2 u_1 \qquad a_2 = u_1$$

$$a_3 = q_3 u_2$$

$$a_1 = ?$$



Need switch between A^+ and A^- on a cell face basis depending on sign of q

Build pos. & neg. flux vector

$$q_n = \min(q(1:N_x), 0)$$

$$q_p = \max(q(z:N_x+1), 0)$$

↑
 N_x+1

