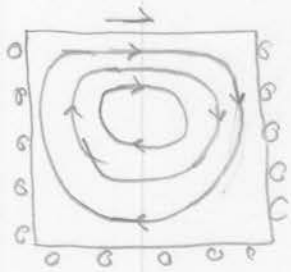


Streamlines & Stream function

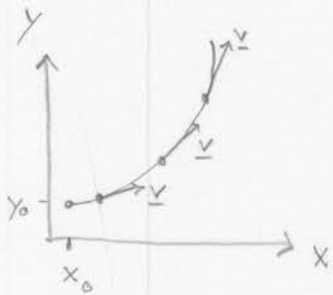
①



Streamlines provide one of the best ways to illustrate flow fields, if applicable.

Definition:

Streamlines are the family of curves that are instantaneously tangent to the velocity field.

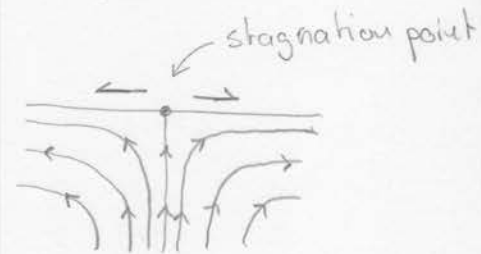


The definition of velocity provides a system of ODE's to compute streamlines:

$$\left. \begin{array}{l} 1) \frac{dx}{dt} = v_x(x) \\ 2) \frac{dy}{dt} = v_y(y) \end{array} \right\} \frac{dy}{dx} = \frac{v_y}{v_x} \quad \text{where } \underline{v}(x) = \begin{pmatrix} v_x(x) \\ v_y(x) \end{pmatrix}$$

Notes: • Safer to solve the system of ODE's because $\frac{dy}{dx}$ may not be bounded and $y(x)$ may be multivalued, as in our example.

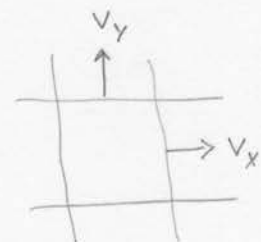
• But the ODE's system has problems with stagnation points because $|v| \rightarrow 0$ and $t \rightarrow \infty$



• Can only determine stagnation points by trial & error.

In Matlab the function `streamline.m` solves for a streamline given an initial point.

Note: `streamline.m` needs all velocities at the same location, so you have to interpolate \underline{v} to the cell centers.



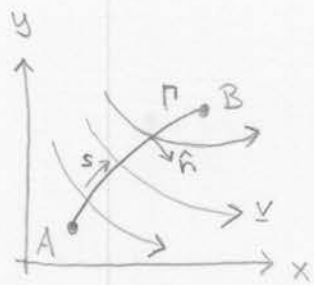
\Rightarrow introduce errors: streamlines hitting no flow bnd's

Different way of thinking about streamlines. (2)

Compute the cumulative flux between A & B.

$$\Psi = \int_{\Gamma} \underline{v} \cdot \hat{n} ds$$

Γ is the path
 s is an arclength variable along Γ
 \hat{n} is the right hand normal



In the absence of fluid sources Ψ should not depend on path. This will be confirmed below.

\Rightarrow Choose a path that simplifies integration.
 We break Γ into two parts along coord. directions.

along Γ_1 : $\underline{v} \cdot \hat{n}_1 = -v_y$ (because \hat{n} points in $-y$ dir.)

along Γ_2 : $\underline{v} \cdot \hat{n}_2 = v_x$

Hence we write integral as: $\Psi = \underbrace{\int_{x_A}^{x_B} -v_y(x, y_A) dx}_{\Gamma_1} + \underbrace{\int_{y_A}^{y_B} v_x(x_B, y) dy}_{\Gamma_2}$

Suppose: $y_A = y_B$



$$\Psi = \int_{x_A}^{x_B} -v_y dx \stackrel{\text{F.T.C.}}{=} \int_{x_A}^{x_B} \frac{\partial \Psi}{\partial x} dx \Rightarrow \frac{\partial \Psi}{\partial x} = -v_y$$

Suppose: $x_A = x_B$



$$\Psi = \int_{y_A}^{y_B} v_x dy \stackrel{\text{F.T.C.}}{=} \int_{y_A}^{y_B} \frac{\partial \Psi}{\partial y} dy \Rightarrow \frac{\partial \Psi}{\partial y} = v_x$$

Therefore:

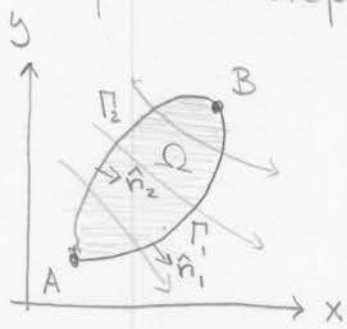
$$\boxed{\frac{\partial \Psi}{\partial x} = -v_y \quad \frac{\partial \Psi}{\partial y} = v_x}$$

This is often given as the definition of the stream function

Physical Interpretation:

- Change of cumulative flux in x -dir is proportional to the negative velocity in y -dir.
- Change of cumulative flux in y -dir is proportional to the velocity in the x -dir.

The conclusions above hold if the integral defining ψ is path independent. (3)



$$\int_{\Gamma_1} \underline{v} \cdot \hat{n}_1 ds = \int_{\Gamma_2} \underline{v} \cdot \hat{n}_2 ds \Rightarrow \int_{\Gamma_1} \underline{v} \cdot \hat{n}_1 ds - \int_{\Gamma_2} \underline{v} \cdot \hat{n}_2 ds = 0$$

Combine paths $\Gamma = \Gamma_1 + \Gamma_2$ and define outside normal, \hat{n} , to the enclosed area Ω .

note that $\hat{n} = \hat{n}_1$ on Γ_1 but $\hat{n} = -\hat{n}_2$ on Γ_2

$$\int_{\Gamma_1} \underline{v} \cdot \hat{n}_1 ds + \int_{\Gamma_2} \underline{v} \cdot \hat{n}_2 ds = \oint_{\Gamma} \underline{v} \cdot \hat{n} ds = 0$$

Hence the Integral is path independent if $\oint_{\Gamma} \underline{v} \cdot \hat{n} ds = 0$

Using the divergence theorem: $\oint_{\Gamma} \underline{v} \cdot \hat{n} ds = \int_{\Omega} \nabla \cdot \underline{v} dV = 0$

Hence the stream function is well defined, if $\boxed{\nabla \cdot \underline{v} = 0}$

- flow is incompressible
 - no sources/sinks of mass
- } not problem for us

Note: In 3D there are two stream functions.

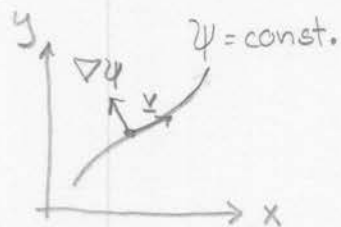
In an incompressible flow without sources the cumulative flux ψ is a unique function of \underline{x} and called the stream function.

What is the relation between ψ and streamlines?

Relation between ψ and streamlines

(4)

1) The level sets of ψ are tangential to the velocity vector.

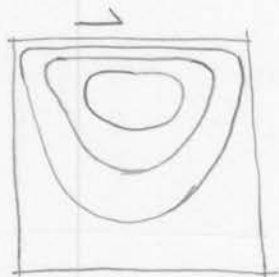


$$\nabla\psi \cdot \underline{v} = \left(\frac{\partial\psi}{\partial x}, \frac{\partial\psi}{\partial y}\right) \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} = (-v_y, v_x) \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} = -v_y v_x + v_x v_y = 0$$

$$\boxed{\nabla\psi \cdot \underline{v} = 0} \Rightarrow \nabla\psi \perp \underline{v}$$

Level sets of ψ are the streamlines.

2)



The magnitude of the velocity is equal to the magnitude of $\nabla\psi$.

$$|\nabla\psi| = \sqrt{(-v_y)^2 + v_x^2} = \sqrt{v_x^2 + v_y^2} = |\underline{v}|$$

If we plot equally spaced contours of ψ the spacing indicates the velocity.

This is the most useful aspect of properly plotted streamlines! Not possible without stream function.

Other applications:

We can compute a ψ for any solenoidal vector field ($\nabla \cdot \underline{g} = 0$)

For example we have used it to visualize the conductive heat flux near a conductive inclusion. However, heat production by decay heating prevents computation of ψ in many cases.