

## Lecture 17: Diffusion

Logistics: - HW3 Thursday is last chance!

- HW6 due Thursday

Last time: - Solute mass balance

$$\phi \frac{\partial c}{\partial t} + \nabla \cdot [q_c - D_H \nabla c] = -\phi k c$$

Advection:  $q_c$

Diffusive:  $-D_H \nabla c$  ind. Diff. + Dispersion

Reactive:  $-\phi k c$

Scaling

$$\frac{\partial c'}{\partial t'} + \nabla' \cdot \left[ \underbrace{\frac{t_c}{E_A}}_{\Pi_1} q' c' - \underbrace{\frac{t_c}{E_D} D'_H \nabla' c'}_{\Pi_2} \right] = - \underbrace{\frac{t_c}{E_R} c'}_{\Pi_3}$$

- Characteristic time scales:

advection:  $t_A = \frac{x_c}{q_c/\phi} = \frac{x_c}{V_c}$

diffusive:  $t_D = \frac{\phi x_c^2}{D_c}$

reactive:  $t_R = \frac{1}{k}$

- Peclet number:  $Pe = \frac{t_P}{t_A} = \frac{q_c x_c}{D_c}$

Today: Diffusion

$$D_a = \frac{t_F}{t_R}$$

# Diffusion

$$\frac{\partial c'}{\partial t'} + \nabla' \cdot \left[ \underbrace{\frac{t_D}{t_A} q' c'}_{Pe} - \frac{t_R}{t_D} D'_H \nabla' c' \right] = - \frac{t_D}{t_R} c'$$

$\frac{t_D}{t_A} \frac{E_A}{t_R} = Pe Da$

choose diff. time scale:  $t_c = t_D$

$$\frac{\partial c'}{\partial t'} + \nabla' \cdot [Pe q' c' - D'_H \nabla' c'] = - Pe Da c'$$

if  $Pe \ll 1$  and  $Da \sim 1$ ,  $\nabla^2 c$

limit equation,

$$\boxed{\frac{\partial c'}{\partial t'} - \nabla' \cdot [D'_H \nabla' c'] = 0}$$

dim. less diffusion equation

re dimensionalize:  $c' = \frac{c}{C_e} \dots$

$$\phi \frac{\partial c}{\partial t} - \nabla \cdot [D_H \nabla c] = 0$$

↑                  ↗

note: if  $q=0$   $D_H = 0$   $D_H = \phi D_w I + D_H^*$

$$\Rightarrow \boxed{\frac{\partial c}{\partial t} - \nabla \cdot [D_w \nabla c] = 0}$$

diff. eqn.

similar eqns:

heat transfer:  $\underbrace{\rho c_p \frac{\partial T}{\partial t}}_{\text{heat storage}} - \nabla \cdot [\kappa \nabla T] = 0$

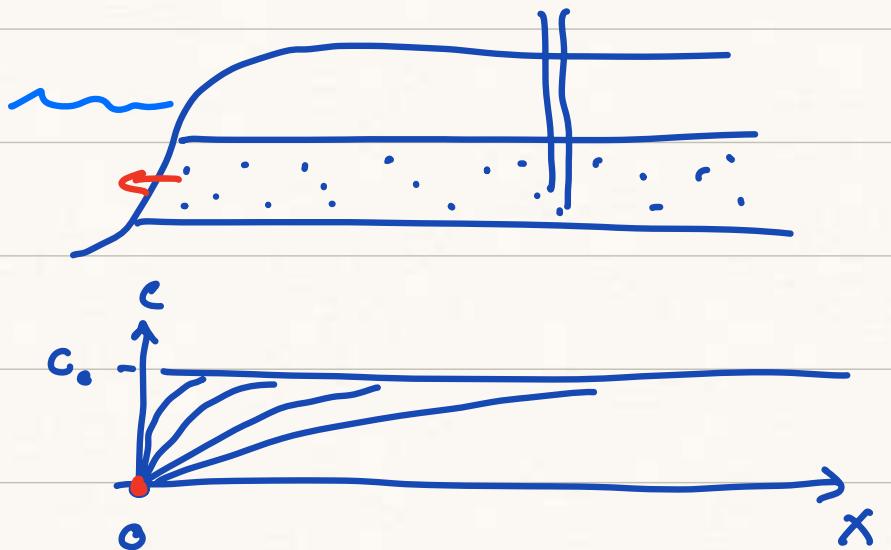
slightly compressible flow:  $\underbrace{S \frac{\partial h}{\partial t}}_{\text{storage}} - \nabla \cdot [\kappa \nabla h] = 0$  incompressible flow

## Diffusion of solute out of aquifer

PDE:  $\frac{\partial c}{\partial t} - D_m \frac{\partial^2 c}{\partial x^2} = 0 \quad x \in [0, \infty)$

BC:  $c(0, t) = 0$

IC:  $c(x, 0) = c_0$



Q: How fast does this concentration front propagate?

How do we scale problem?

No obvious  $\uparrow$  length scale, because we have half-space exterior

Internal length scale:

$$\sqrt{D_m E} \quad \left\{ \sqrt{\frac{m^2}{s}} s = m \right.$$

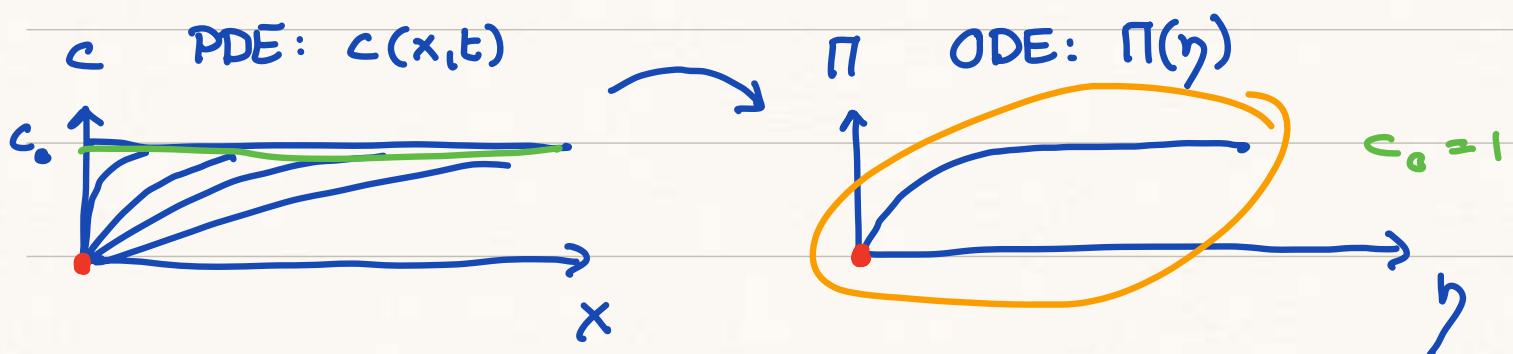
What about  $x' = \frac{x}{\sqrt{D_m t}}$  ?

We are not just scaling with parameters, but with another independent variable

⇒ results in a new independent variable !

$$\boxed{\gamma = \frac{x}{\sqrt{4Dt}}}$$

Boltzmann variable



By introducing  $\gamma \sim \frac{x}{\sqrt{t}}$  we reduce PDE to an ODE

$\gamma$  is called the similarity variable

and solution is said to be self-similar

Q: What is the ODE ?

First :  $c' = \frac{c}{c_0}$        $0 \leq c' \leq 1$       IC:  $c' = 1$

Solution :  $c'(x,t) = \underline{\Pi(\gamma(x,t))}$        $\Pi(\gamma)$

substitute into PDE

Transform derivatives:

$$\frac{\partial c'}{\partial t} = \frac{\partial}{\partial t}(\Pi(y(x,t))) = \frac{d\Pi}{dy} \frac{\partial y}{\partial t} \quad \text{chain rule}$$
$$y = \frac{x}{\sqrt{4Dt}}$$

$$\frac{\partial y}{\partial x} = \frac{1}{\sqrt{4Dt}} \quad \frac{\partial y}{\partial t} = -\frac{1}{2t}$$

$$\frac{\partial^2 c'}{\partial x^2} = \frac{\partial}{\partial x}\left(\frac{d\Pi}{dy} \frac{\partial y}{\partial x}\right) = \left(\frac{\partial y}{\partial x}\right)^2 \frac{d^2\Pi}{dy^2} = \frac{1}{4Dt} \frac{d^2\Pi}{dy^2}$$

Substitute into PDE:

$$\frac{\partial c'}{\partial t} - D_m \frac{\partial^2 c'}{\partial x^2} = -\frac{1}{2t} \frac{d\Pi}{dy} - \frac{Dm}{4Dt} \frac{d^2\Pi}{dy^2} = 0 \quad \Leftarrow$$

ODE:

$$\frac{d^2\Pi}{dy^2} + 2y \frac{d\Pi}{dy} = 0$$

BC:

$$\Pi(0) = 0$$

$$\lim_{y \rightarrow \infty} \Pi = 1$$

Self-Similar ODE

$$y = \frac{x}{\sqrt{4Dt}}$$

↑  
IC of PDE     $t \rightarrow 0$      $y \rightarrow \infty$   
( $x$  finite)

Solve ODE:

1) substitute:  $u = \frac{d\Pi}{dy} \Rightarrow \frac{du}{dy} + z\gamma u = 0$

2) separate variables:  $\frac{du}{u} = -z\gamma dy$   
 $\ln u = -z^2 + a$   
 $u = b e^{-z^2}$

3) re substitute:  $\frac{d\Pi}{dy} = b e^{-z^2}$

4) separate variables:  $d\Pi = b e^{-z^2} dy$   
 $\Pi(y) = b \int_0^y e^{-z^2} dz$

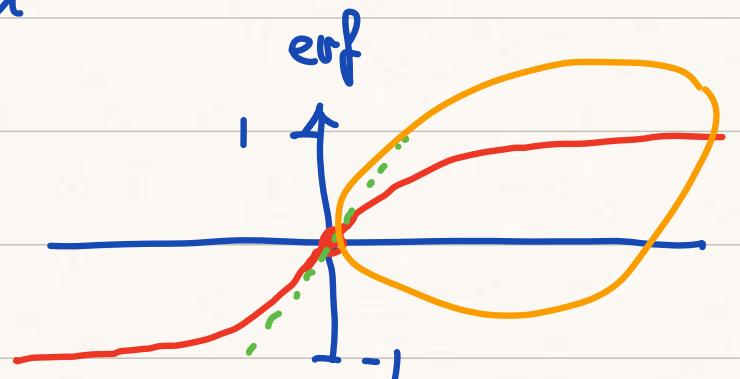
$z = \text{dummy variable}$

integral does not have analytic solu.

$\Rightarrow$  give it a name and move on

5) identify error function

$$\boxed{\text{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-z^2} dz}$$



Properties:

- $\text{erf}(-x) = -\text{erf}(x)$  "point symmetric"

- $\operatorname{erf}(0) = 0$
- $\operatorname{erf}(x) \approx x$        $|x| \ll 1$
- $\lim_{x \rightarrow \infty} \operatorname{erf}(x) = 1$ ,     $\lim_{x \rightarrow -\infty} \operatorname{erf}(x) = -1$

Therefore:  $\Pi(y) = b \frac{\sqrt{\pi}}{2} \operatorname{erf}(y)$

$$\text{BC: } \lim_{y \rightarrow \infty} \Pi(y) = b \frac{\sqrt{\pi}}{2} = 1 \Rightarrow b = \frac{2}{\sqrt{\pi}}$$

Self-similar solu:

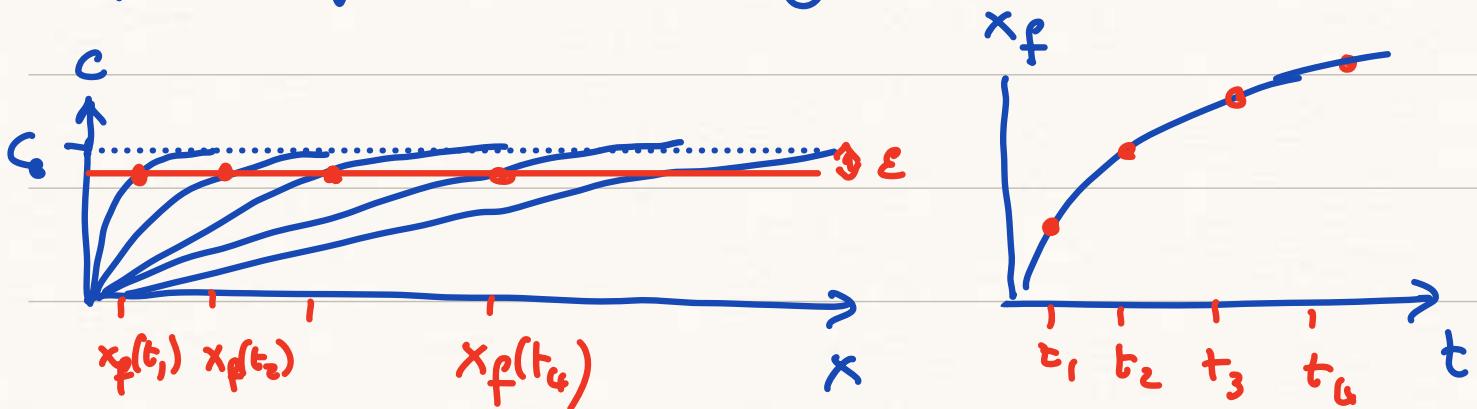
$$\boxed{\Pi(y) = \operatorname{erf}(y)}$$

6) Resubstitute:  $c = c_0 c' \quad y = \frac{x}{\sqrt{4DE}}$

Transient concentration evolution

$$\boxed{c(x,t) = c_0 \operatorname{erf}\left(\frac{x}{\sqrt{4DE}}\right)}$$

# Speed of front propagation



Note: infinitesimal changes propagate infinitely fast!

Finite changes,  $\Delta c = \varepsilon \geq 1$ , propagate with finite velocity.

Define front as location where  $c$  has changed by  $\varepsilon$  from its initial value  $c_0$

$$c(x_f, t) = c_0 - \varepsilon c_0 = c_0 (1 - \varepsilon)$$

$$\frac{c_0}{\varepsilon} \operatorname{erf}\left(\frac{x_f}{\sqrt{4Dt}}\right) = 1 - \varepsilon$$

inverse error function

$$\frac{x_f}{\sqrt{4Dt}} = \operatorname{erf}^{-1}(1 - \varepsilon) = \alpha(\varepsilon) = \text{const}$$

$\Rightarrow$

$$x_f = \alpha(\epsilon) \sqrt{4Dt}$$

Concentration perturbation propagated

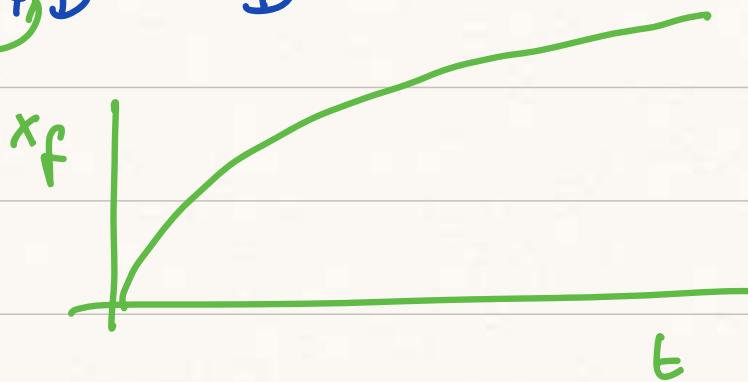
as  $\sqrt{t}$

$$x_f \sim \sqrt{t}$$

This matches our diffusive time scale

$$x_c = x_f = \alpha(\epsilon) \sqrt{4Dt_D}$$

$$t_D = \frac{x_c^2}{\alpha(\epsilon) 4D} \sim \frac{x_c^2}{D}$$



$$\frac{\partial c}{\partial t} - \underbrace{\nabla \cdot [D_m \nabla c]}_{\text{known}} = 0$$

new